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Anexo A

Método simplificado para avaliação do potencial da liquefação de solos.

Um método simplificado para determinar o potencial da liquefação de solos, proposto por Seed e Idriss [Seed, H.B. & Idriss, I.M., 1971] e aperfeiçoado por diversos autores ([Finn, W.D.L., 1993], [Youd, T.L., et al., 2001], [Seed, R.B., et al., 2003]) consiste em verificar em determinada profundidade se as tensões geradas pelo carregamento sísmico excedem à resistência ao cisalhamento do solo na condição residual.

Este critério de análise permite estabelecer um fator de segurança contra a liquefação, FS, definido por

$$FS = \frac{CRR}{CSR} \quad (\text{Eq. A.1})$$

onde CRR é a razão de resistência cíclica na liquefação (*Cyclic Resistance Ratio*) e CSR a razão de tensão cíclica gerado pelo terremoto de projeto (*Cyclic Stress Ratio*). Se o fator de segurança for menor que a unidade ($FS < 1$), então a liquefação deve ocorrer.

A razão de resistência cíclica, CRR, é determinada da figura A.1 [Seed, R.B., et al., 2003], onde o número de golpes corrigidos do ensaio SPT, $(N_1)_{60}$, é estimado pela relação

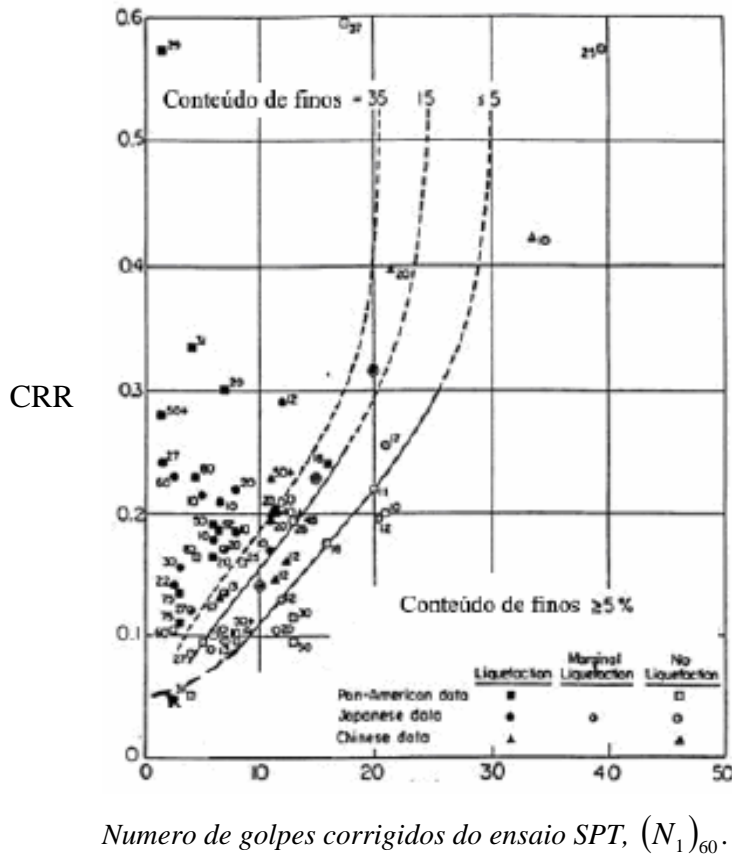
$$(N_1)_{60} = (N)_{60} \cdot C_N \quad (\text{Eq. A.2})$$

onde C_N é um fator de correção e $(N)_{60}$ o número de golpes do ensaio SPT.

Liao e Whitman [Liao, S.S.C. & Whitman, R.V., 1986] sugerem a seguinte equação para este fator de correção,

$$C_N = \frac{1}{\sqrt{\sigma'_{v0}}} \quad (\text{Eq. A.3})$$

onde σ'_{v0} é a tensão vertical efetiva inicial expressa em kgf/cm^2 . Outras correlações similares para cálculo de C_N podem ser encontradas na literatura ([Youd, T.L., et al., 2001], [Seed, R.B., et al., 2003]).



Numero de golpes corrigidos do ensaio SPT, $(N_1)_{60}$.

Figura A.1 Razão de resistência cíclica (CRR) versus número de golpes corrigidos do ensaio SPT, $(N_1)_{60}$, para terremotos com magnitude igual a 7,5 [Youd, T.L., et al., 2001].

A razão de tensão cíclica, CSR, é expressa por

$$CSR = CSR_{eq, Mw=7,5} / DWF_M / K_\sigma \quad (\text{Eq. A.4})$$

onde $CSR_{eq, Mw=7,5}$ é a razão de tensão cíclica equivalente correspondente a um terremoto de magnitude (M_w) igual a 7,5; DWF_M o fator de correção da magnitude do terremoto (*Magnitude-Related Duration Weighting Factor*) e K_σ o fator de correção da tensão vertical efetiva. As figuras A.2 e A.3 são utilizadas para a determinação de DWF_M e K_σ , respectivamente. Kramer [Kramer, S.L., 1996] recomenda também a utilização da tabela A.1 para a estimativa do fator de correção da magnitude do sismo.

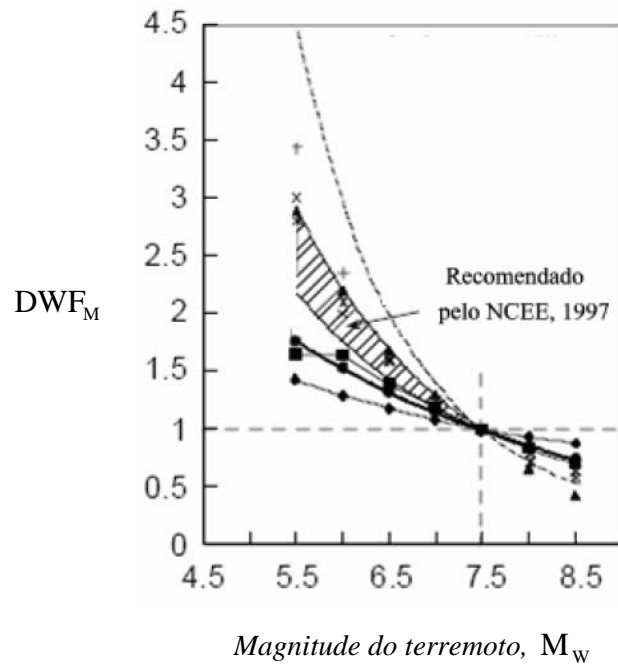


Figura A.2 Fator de correção da magnitude do terremoto [Seed, R.B., et al., 2003].

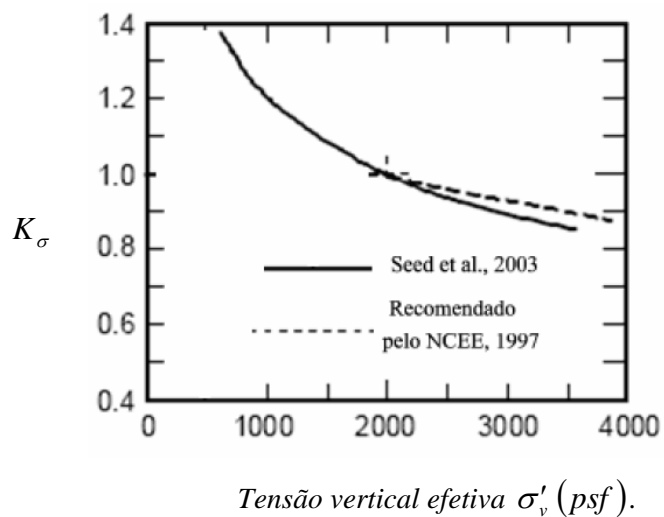


Figura A.3 Fator de correção devido à tensão inicial de cisalhamento [Seed, R.B., et al., 2003].

A razão de tensão cíclica equivalente com $M_w = 7,5$, $CSR_{eq Mw=7,5}$, pode ser escrita como

$$CSR_{eq Mw=7,5} = 0,65 CSR_{max Mw=7,5} \quad (\text{Eq. A.5})$$

com

$$CSR_{eq Mw=7,5} = \frac{a_{max}}{\bar{g}} \left(\frac{\sigma_v}{\sigma'_{vo}} \right) r_d \quad (\text{Eq. A.6})$$

onde $CSR_{\max, Mw=7,5}$ é a razão de tensão cíclica máxima com $M_w = 7,5$; a_{\max} a aceleração horizontal máxima do terremoto na superfície; \bar{g} a aceleração da gravidade; r_d um fator de redução da tensão cíclica devido à profundidade; σ_v a tensão total vertical e σ'_{v0} a tensão vertical efetiva inicial.

De acordo com Finn [Finn, W.D.L., 1993], na prática japonesa o fator de redução de tensão devido à profundidade é frequentemente aproximado pela seguinte correlação

$$r_d = 1 - 0,0015 z \quad (\text{Eq. A.7})$$

onde z indica a profundidade do terreno em metros. Youd e colaboradores [Youd, T.L., et al., 2001] recomendam a utilização da figura A.4 para determinação do fator de redução da tensão cíclica devido à profundidade.

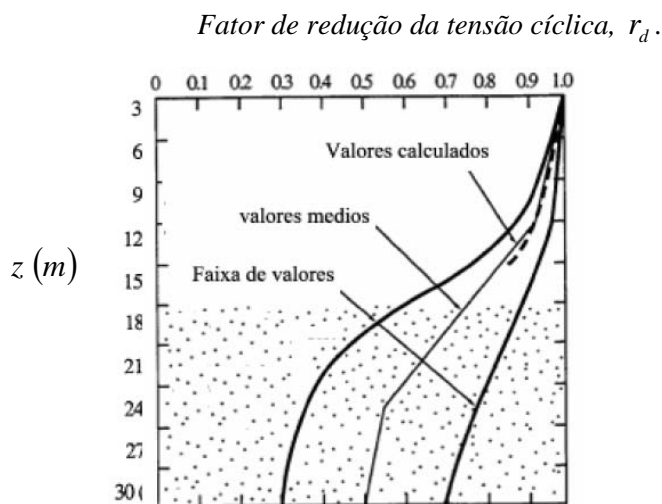


Figura A.4 Fator de redução da tensão cíclica devido à profundidade [Youd, T.L., et al., 2001].

Tabela A.1 Fator de correção da magnitude do sismo [Kramer, S.L., 1996].

Magnitude do terremoto, M_w	Fator de correção, DWF_M
$5\frac{1}{4}$	1,50
5	1,32
$6\frac{3}{4}$	1,13
$7\frac{1}{2}$	1,00
$8\frac{1}{2}$	0,89

Anexo B

Discretização das equações fundamentais para a condição não-saturada na forma u-p.

As equações fundamentais em meios porosos saturados na forma **u-w-p** são [Zienkiewicz, O.C., et al., 1999]:

$$\sigma_{ij,j} - \rho \ddot{u}_i - \rho_w \left(\ddot{w}_i + \dot{w}_j \dot{w}_{ij} \right) + \rho b_i = 0 \quad (\text{Eq. B.1})$$

$$-p_{w,i} - R_i - \rho_w \ddot{u}_i - \frac{\rho_w}{n} \left(\ddot{w}_i + \dot{w}_j \dot{w}_{i,j} \right) + \rho_w b_i = 0 \quad (\text{Eq. B.2})$$

$$\dot{w}_{i,i} + \dot{\varepsilon}_{ii} + \frac{n\dot{p}_w}{K_w} + \frac{(1-n)\dot{p}_w}{K_s} - \frac{K_T}{K_s} \left(\dot{\varepsilon}_{ii} + \frac{\dot{p}_w}{K_s} \right) + n \frac{\dot{\rho}_w}{\rho_w} + \dot{s}_0 = 0 \quad (\text{Eq. B.3})$$

onde as equações B.1 e B.2 representam as equações de movimento para o sistema sólido-fluido e para o fluido respectivamente, e a equação B.3 representa a equação de continuidade do fluido.

Com o objeto de estabelecer as equações fundamentais para descrever o comportamento dinâmico do meio poroso em condição não-saturada (sistema sólido-água-ar) algumas modificações devem ser introduzidas:

- ✓ Os vazios contidos no meio poroso são preenchidos parcialmente por água e parcialmente por ar, resultando

$$S_{r_w} + S_{r_a} = 1 \quad (\text{Eq. B.4})$$

onde S_{r_w} e S_{r_a} é o grau de saturação da água e do ar, respectivamente.

- ✓ A densidade do meio poroso é expressa por

$$\rho = nS_{r_w}\rho_w + nS_{r_a}\rho_a + (1-n)\rho_s \quad (\text{Eq. B.5})$$

onde ρ_w , ρ_a e ρ_s são as massas específicas da água, do ar e do sólido respectivamente.

- ✓ A pressão do ar é desconsiderada, i.e. $P_a = 0$ [Zienkiewicz, O.C., et al., 1990b].
- ✓ O princípio das tensões efetivas [Terzaghi, K., 1936] pode ser adaptado para solos parcialmente saturados modificando-se a expressão da poropressão de acordo com [Bishop, A.W. & Blight, G.E., 1963],

$$P_{av} = \chi p_w + (1 - \chi) P_a \quad (\text{Eq. B.6})$$

onde χ é um parâmetro que depende do grau de saturação do sistema. Uma boa aproximação de χ pode ser dada pelo grau de saturação da água, S_{rw} , [Zienkiewicz, O.C., et al., 1999],

$$\chi = S_{rw} \quad (\text{Eq. B.7})$$

Desconsiderando-se o valor da pressão do ar ($P_a = 0$), a equação B.6 pode ser aproximada por,

$$p_{av} = S_{rw} p_w \quad (\text{Eq. B.8})$$

modificando-se o princípio das tensões efetivas para solos saturados para:

$$\sigma'_{ij} = \sigma_{ij} + \delta_{ij} S_{rw} p_w \quad (\text{Eq. B.9})$$

- ✓ Considerando a permeabilidade dependente do grau de saturação, vem

$$k = k \langle S_{rw} \rangle \quad (\text{Eq. B.10})$$

Vários estudos reportados na literatura estabelecem relações entre o grau de saturação, permeabilidade e a poropressão ([Huang, M. & Zienkiewicz O.C., 1998], [Alonso, E.E., et al., 1987], [Bear, J., et al., 1984], [Lloret, A. & Alonso, E.E., 1980], [Safai, N.M. & Pinder, G.F., 1979], [Narasimhan, T.N. & Witherspoon, P.A., 1978], [Van Genuchten, M.T.; Pinder, G.F. & Saukin, W.P., 1977], [Neuman, S.P., 1975], [Liakopoulos, A.C., 1965]).

Levando-se em conta as observações feitas anteriormente para solos não-saturados, as equações B.1, B.2 e B.3 podem então ser modificadas.

A equação de continuidade do fluido (equação B.3) é modificada para,

$$\begin{aligned} \dot{w}_{i,i} + \dot{\varepsilon}_{ii} + \frac{n}{K_w} S_{rw} \dot{p}_w + \frac{(1-n)}{K_s} S_{rw} \dot{p}_w - \frac{K_T}{K_s} \left(\dot{\varepsilon}_{ii} + \frac{1}{K_s} S_{rw} \dot{p}_w \right) \\ + n \frac{\dot{p}_w}{\rho_w} + \dot{s}_0 = 0 \end{aligned} \quad (\text{Eq. B.11})$$

ou

$$\dot{w}_{i,i} + \tilde{\alpha} \dot{\varepsilon}_{ii} + \frac{\dot{p}_w}{\tilde{Q}^*} + n S_{rw} \frac{\dot{p}_w}{\rho_w} + \dot{s}_0 = 0 \quad (\text{Eq. B.12})$$

onde

$$\frac{1}{\tilde{Q}^*} \equiv S_{rw} \left(\frac{n}{K_w} + \frac{\tilde{\alpha} - n}{K_s} \right) + n \frac{\dot{S}_{rw}}{\dot{p}_w} \quad (\text{Eq. B.13})$$

sendo

$$\tilde{\alpha} = 1 - \frac{K_T}{K_s} \quad (\text{Eq. B.14})$$

A definição da compressibilidade equivalente do sistema sólido-agua-ar (condição não-saturada) pode ser expressa como

$$\tilde{C}_T^* \equiv S_{rw} [nC_w + (\tilde{\alpha} - n)C_s] + n \frac{\dot{S}_{rw}}{\dot{p}_w} \quad (\text{Eq. B.15})$$

onde $\tilde{C}_T^* = 1/\tilde{Q}^*$ é a compressibilidade equivalente do sistema sólido-agua-ar; $C_w = 1/K_w$ a compressibilidade do fluido (água) e $C_s = 1/K_s$ a compressibilidade do sólido. A equação B.15 também pode ser escrita em termos da compressibilidade equivalente do sistema sólido-agua (condição saturada) C_T .

$$\tilde{C}_T^* \equiv S_{rw} C_T + n \frac{\dot{S}_{rw}}{\dot{p}_w} \quad (\text{Eq. B.16})$$

com

$$C_T = nC_w + (\tilde{\alpha} - n)C_s \quad (\text{Eq. B.17})$$

No caso das equações B.1 e B.2 a componente de tensão total deve ser substituída em função do princípio das tensões efetivas na condição não-saturada (equação B.9).

A formulação simplificada **u-p**, com a eliminação da variável \dot{w}_i nas equações acima, pode então ser escrita, na sua forma incremental, para meios porosos não-saturados,

$$\delta \sigma_{ij,j} - \rho \delta \ddot{u}_i + \rho \delta b_i = 0 \quad (\text{Eq. B.18})$$

$$\left(k_{ij}(-\delta p_{w,j} - S_{r,w}\rho_w \delta \ddot{u}_j + S_{r,w}\rho_w \delta b_j)\right)_j + \delta \ddot{u}_{i,i} + \frac{\delta \dot{p}_w}{\tilde{Q}^*} = 0 \quad (\text{Eq. B.19})$$

com

$$\delta \sigma_{ij} = \delta \sigma'_{ij} - \delta_{ij} S_{r,w} \delta p_w \quad (\text{Eq. B.20})$$

Com o objetivo de obter a solução numérica das equações governantes, é necessário discretizar estas equações, tanto espacial quanto temporalmente, conservando como variáveis primárias os incrementos de deslocamento nodal do sólido e da poropressão nodal do fluido.

Aplicando-se o método de Galerkin na equação B.18, obtém-se a seguinte equação discretizada a nível local para o sólido:

$$\mathbf{M} \delta \ddot{\mathbf{u}} + \mathbf{P} \langle \delta \bar{\mathbf{u}} \rangle - \mathbf{Q} \delta \bar{\mathbf{p}}_w - \delta \bar{\mathbf{f}}^{(s)} = \mathbf{0} \quad (\text{Eq. B.21})$$

com

$$\mathbf{M} = \int_{\Omega} \mathbf{N}^{uT} \rho \mathbf{N}^u d\Omega \quad (\text{Eq. B.22})$$

$$\mathbf{P} \langle \delta \bar{\mathbf{u}} \rangle = \int_{\Omega} \mathbf{B}^{uT} \delta \boldsymbol{\sigma} d\Omega \quad (\text{Eq. B.23})$$

$$\mathbf{Q} = \int_{\Omega} \mathbf{B}^{uT} S_{r,w} \mathbf{m} \mathbf{N}^w d\Omega \quad (\text{Eq. B.24})$$

$$\delta \bar{\mathbf{f}}^{(s)} = \int_{\Omega} \mathbf{N}^{uT} \mathbf{b} d\Omega + \int_{\Gamma} \mathbf{N}^{uT} \mathbf{t} d\Gamma \quad (\text{Eq. B.25})$$

De forma similar, considerando-se a equação B.19 resulta a seguinte equação discretizada a nível local para o fluido:

$$\mathbf{Q}^T \delta \dot{\mathbf{u}} + \mathbf{H} \delta \bar{\mathbf{p}}_w + \mathbf{S} \delta \dot{\bar{\mathbf{p}}}_w - \delta \bar{\mathbf{f}}^{(w)} = \mathbf{0} \quad (\text{Eq. B.26})$$

com

$$\mathbf{H} = \int_{\Omega} \nabla \mathbf{N}^{wT} \mathbf{k} \nabla \mathbf{N}^w d\Omega \quad (\text{Eq. B.27})$$

$$\mathbf{S} = \int_{\Omega} \mathbf{N}^{wT} \frac{1}{\tilde{Q}^*} \mathbf{N}^w d\Omega \quad (\text{Eq. B.28})$$

$$\delta \bar{\mathbf{f}}^{(w)} = - \int_{\Omega} \mathbf{N}^{wT} \nabla^T (\mathbf{k} S_{r,w} \rho_w \bar{\mathbf{b}}) d\Omega + \int_{\Gamma} \mathbf{N}^{wT} \mathbf{q} d\Gamma \quad (\text{Eq. B.29})$$

Combinando as equações B.21 e B.26 para descrever, a nível local, o comportamento dinâmico do elemento acoplado (sólido-água-ar), tem-se

$$\begin{bmatrix} \mathbf{M} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{Bmatrix} \delta \ddot{\mathbf{u}} \\ \delta \ddot{\mathbf{p}}_w \end{Bmatrix} + \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{Q}^T & \mathbf{S} \end{bmatrix} \begin{Bmatrix} \delta \dot{\mathbf{u}} \\ \delta \dot{\mathbf{p}}_w \end{Bmatrix} + \begin{bmatrix} \mathbf{K} & -\mathbf{Q} \\ \mathbf{0} & \mathbf{H} \end{bmatrix} \begin{Bmatrix} \delta \bar{\mathbf{u}} \\ \delta \bar{\mathbf{p}}_w \end{Bmatrix} - \begin{Bmatrix} \delta \bar{\mathbf{f}}^{(s)} \\ \delta \bar{\mathbf{f}}^{(w)} \end{Bmatrix} = \mathbf{0} \quad (\text{Eq. B.30})$$

Após o procedimento de montagem dos elementos finitos acoplados, a equação discreta que descreve o comportamento dinâmico acoplado sólido-água-ar a nível global (sistema), tem a seguinte forma,

$$\begin{bmatrix} \tilde{\mathbf{M}} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{Bmatrix} \delta \ddot{\mathbf{u}} \\ \delta \ddot{\mathbf{p}}_w \end{Bmatrix} + \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \tilde{\mathbf{Q}}^T & \tilde{\mathbf{S}} \end{bmatrix} \begin{Bmatrix} \delta \dot{\mathbf{u}} \\ \delta \dot{\mathbf{p}}_w \end{Bmatrix} + \begin{bmatrix} \tilde{\mathbf{K}} & -\tilde{\mathbf{Q}} \\ \mathbf{0} & \tilde{\mathbf{H}} \end{bmatrix} \begin{Bmatrix} \delta \tilde{\mathbf{u}} \\ \delta \tilde{\mathbf{p}}_w \end{Bmatrix} - \begin{Bmatrix} \delta \tilde{\mathbf{f}}^{(s)} \\ \delta \tilde{\mathbf{f}}^{(w)} \end{Bmatrix} = \mathbf{0} \quad (\text{Eq. B.31})$$

O processo de solução numérica se completa com a discretização temporal das componentes da *equação de equilíbrio dinâmico do sistema* (equação B.31) no tempo $t + \Delta t$,

$$\tilde{\mathbf{M}} \delta \ddot{\mathbf{u}}_{t+\Delta t} + \tilde{\mathbf{P}} \langle \delta \tilde{\mathbf{u}}_{t+\Delta t} \rangle - \tilde{\mathbf{Q}} \delta \tilde{\mathbf{p}}_{wt+\Delta t} - \delta \tilde{\mathbf{f}}^{(s)} = \mathbf{0} \quad (\text{Eq. B.32})$$

$$\tilde{\mathbf{Q}} \delta \dot{\mathbf{u}}_{t+\Delta t} + \tilde{\mathbf{H}} \delta \tilde{\mathbf{p}}_{wt+\Delta t} + \tilde{\mathbf{S}} \delta \dot{\tilde{\mathbf{p}}}_{wt+\Delta t} - \delta \tilde{\mathbf{f}}^{(w)} = \mathbf{0} \quad (\text{Eq. B.33})$$

através do método de Newmark Generalizado, GN_{ij} .

Considerando-se o esquema GN22 para o sólido,

$$\delta \tilde{\mathbf{u}}_{t+\Delta t} = \delta \tilde{\mathbf{u}}_t + \delta \dot{\tilde{\mathbf{u}}}_t \Delta t + \frac{1}{2} \delta \ddot{\tilde{\mathbf{u}}}_t \Delta t^2 + \frac{1}{2} \beta_2 (\delta \ddot{\tilde{\mathbf{u}}}_{t+\Delta t} - \delta \ddot{\tilde{\mathbf{u}}}_t) \Delta t^2 \quad (\text{Eq. B.34})$$

$$\delta \dot{\tilde{\mathbf{u}}}_{t+\Delta t} = \delta \dot{\tilde{\mathbf{u}}}_t + \delta \ddot{\tilde{\mathbf{u}}}_t \Delta t + \beta_1 (\delta \ddot{\tilde{\mathbf{u}}}_{t+\Delta t} - \delta \ddot{\tilde{\mathbf{u}}}_t) \Delta t \quad (\text{Eq. B.35})$$

$$\delta \ddot{\tilde{\mathbf{u}}}_{t+\Delta t} = \frac{1}{\beta_1 \Delta t^2} (\delta \tilde{\mathbf{u}}_{t+\Delta t} - \delta \tilde{\mathbf{u}}_t) - \frac{1}{\beta_1 \Delta t} \delta \dot{\tilde{\mathbf{u}}}_t - \left(\frac{1}{2\beta_1} - 1 \right) \delta \ddot{\tilde{\mathbf{u}}}_t \quad (\text{Eq. B.36})$$

e GN11 para o fluido,

$$\delta \tilde{\mathbf{p}}_{wt+\Delta t} = \delta \tilde{\mathbf{p}}_{wt} + \delta \dot{\tilde{\mathbf{p}}}_{wt} \Delta t + \bar{\beta}_1 (\delta \tilde{\mathbf{p}}_{wt+\Delta t} - \delta \tilde{\mathbf{p}}_{wt}) \Delta t \quad (\text{Eq. B.37})$$

$$\delta \dot{\tilde{\mathbf{p}}}_{wt+\Delta t} = \delta \dot{\tilde{\mathbf{p}}}_{wt} + \frac{\delta \tilde{\mathbf{p}}_{wt+\Delta t} - \delta \tilde{\mathbf{p}}_{wt}}{\Delta t} \quad (\text{Eq. B.38})$$

onde as variáveis no tempo $t + \Delta t$ são quantidades a calcular e as variáveis no tempo t são valores conhecidos ou previamente determinados.

A equação de *equação de equilíbrio dinâmico do sistema* (equação B.31), pode ser escrita de forma equivalente ao *equilíbrio estático*,

$$\overline{\overline{\mathbf{K}}}_S \delta \tilde{\Phi}_{t+\Delta t} = \overline{\overline{\delta \mathbf{f}}}_{S t+\Delta t} \quad (\text{Eq. B.39})$$

com

$$\overline{\overline{\mathbf{K}}}_S = \begin{bmatrix} \frac{1}{\alpha_s \Delta t^2} \tilde{\mathbf{M}} + \frac{\beta_s}{\alpha_s \Delta t} \tilde{\mathbf{C}} + \tilde{\mathbf{K}} & -\tilde{\mathbf{Q}} \\ \frac{\beta_w}{\alpha_w \Delta t} \tilde{\mathbf{Q}}^T & \frac{\beta_w}{\alpha_w \Delta t} \tilde{\mathbf{S}} + \tilde{\mathbf{H}} \end{bmatrix} \quad (\text{Eq. B.40})$$

$$\delta \tilde{\Phi}_{t+\Delta t} = \left\{ \delta \tilde{\mathbf{u}}_{t+\Delta t} \quad \delta \tilde{\mathbf{p}}_{w t+\Delta t} \right\}^T \quad (\text{Eq. B.41})$$

$$\overline{\overline{\delta \mathbf{f}}}_{S t+\Delta t} = \left\{ \overline{\overline{\delta \mathbf{f}}}_{t+\Delta t}^{(s)} \quad \overline{\overline{\delta \mathbf{f}}}_{t+\Delta t}^{(w)} \right\}^T \quad (\text{Eq. B.42})$$

onde

$$\begin{aligned} \overline{\overline{\delta \mathbf{f}}}_{t+\Delta t}^{(s)} &= -\delta \tilde{\mathbf{f}}_{t+\Delta t}^{(s)} + \delta \tilde{\mathbf{f}}_t^{(s)} \\ &+ \left(\frac{1}{\alpha_s \Delta t} \delta \tilde{\mathbf{u}}_t + \frac{1}{2\alpha_s} \delta \tilde{\mathbf{u}}_t \right) \tilde{\mathbf{M}} \\ &+ \left(\frac{\beta_s}{\alpha_s} \delta \dot{\tilde{\mathbf{u}}}_t + \Delta t \left(\frac{\beta_s}{2\alpha_s} - 1 \right) \delta \ddot{\tilde{\mathbf{u}}}_t \right) \tilde{\mathbf{C}} \end{aligned} \quad (\text{Eq. B.43})$$

$$\begin{aligned} \overline{\overline{\delta \mathbf{f}}}_{t+\Delta t}^{(w)} &= -\delta \tilde{\mathbf{f}}_{t+\Delta t}^{(w)} + \delta \tilde{\mathbf{f}}_t^{(w)} \\ &+ \left(\frac{\beta_s}{\alpha_s} \delta \dot{\tilde{\mathbf{u}}}_t + \Delta t \left(\frac{\beta_s}{2\alpha_s} - 1 \right) \delta \ddot{\tilde{\mathbf{u}}}_t \right) \tilde{\mathbf{Q}}^T \\ &+ \left(\frac{\beta_w}{\alpha_w} \delta \dot{\tilde{\mathbf{p}}}_{w t} \right) \tilde{\mathbf{S}} \end{aligned} \quad (\text{Eq. B.44})$$

A utilização da equação B.39 como procedimento de solução numérica permite a determinação direta das variáveis primárias $\delta \tilde{\mathbf{u}}_{t+\Delta t}$ e $\delta \tilde{\mathbf{p}}_{t+\Delta t}$.

Finalmente, é necessário incorporar na formulação das equações fundamentais discretas relações constitutivas (em termos de tensões efetivas) de modo de obter as equações governantes totalmente discretizadas e prever o comportamento dinâmico de solos não-saturados através da sua solução numérica. Neste sentido, uma adaptação da teoria da plasticidade generalizada feita por

[Bolzon, G., et al., 1996] para simulação de carregamentos sob condições parcialmente saturadas, pode ser utilizada. Para o caso de solos cimentados parcialmente saturados, Yang e colaboradores [Yang, C., et al., 2008] formularam também um modelo constitutivo combinando conceitos da teoria da plasticidade generalizada com o modelo BBB (Modelo Básico Barcelona) proposto por Alonso [Alonso, E.E., et al., 1990].

Anexo C

Diagrama de blocos do programa computacional

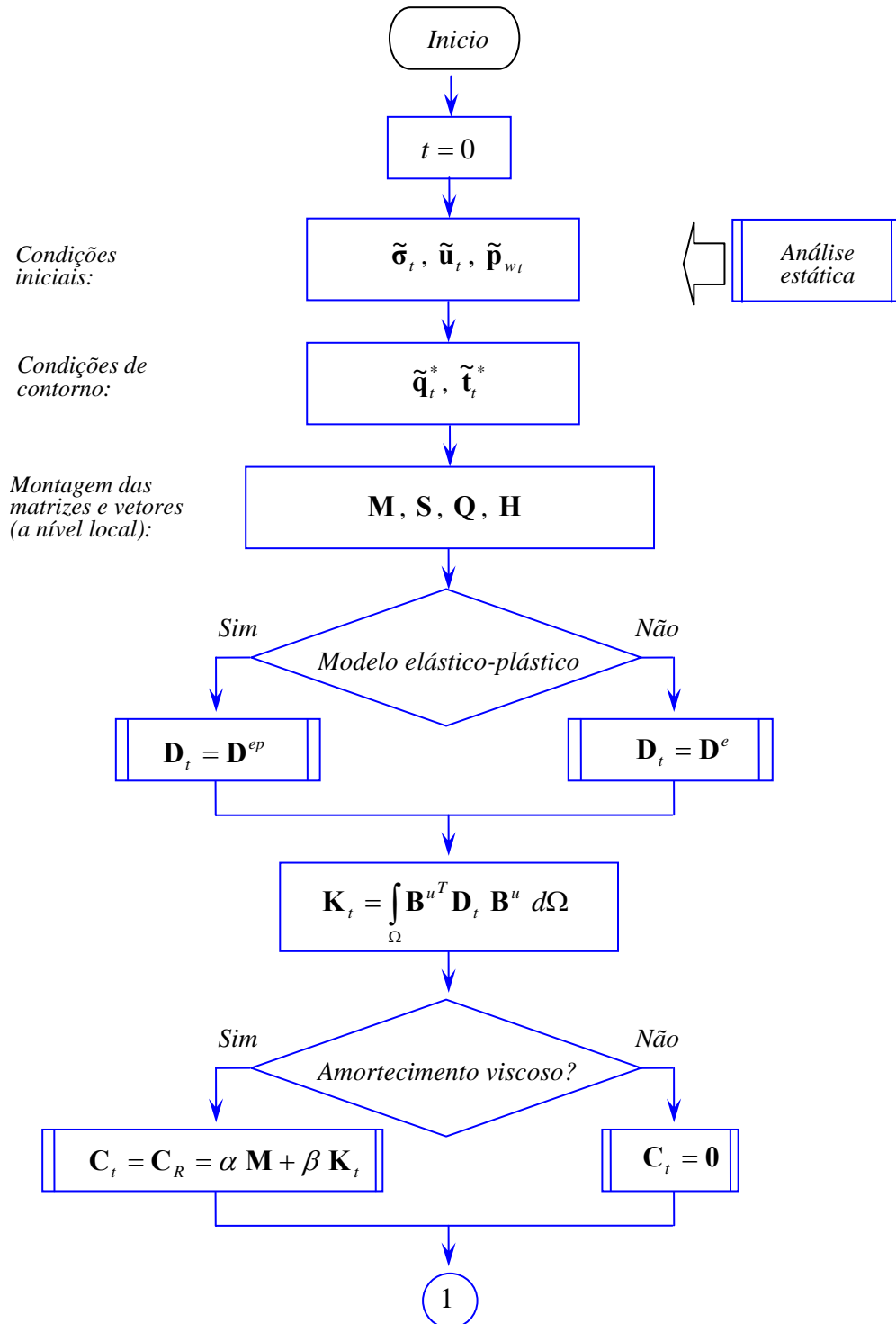


Figura C.1(a) Esquema geral do programa desenvolvido nesta pesquisa.

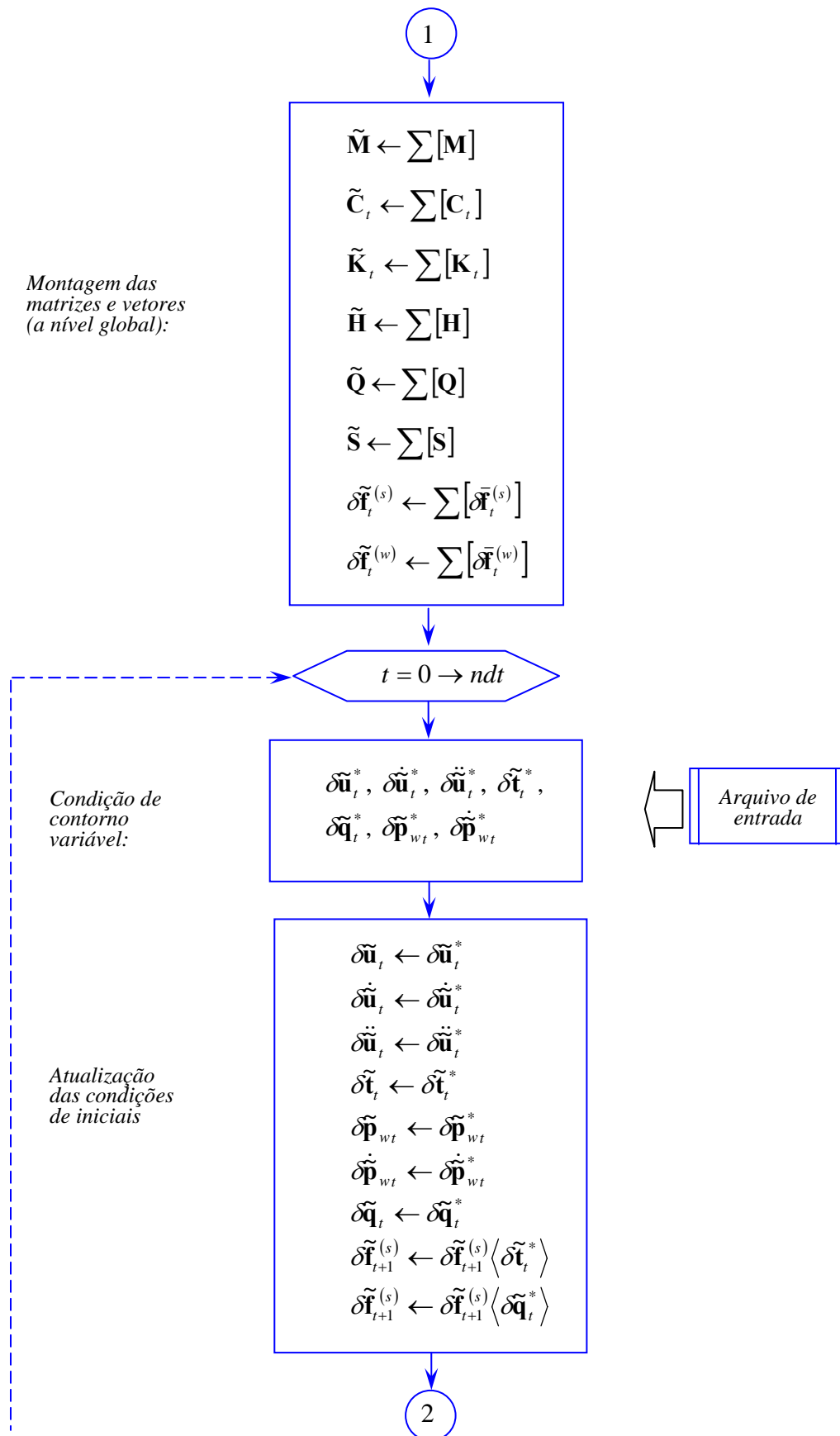


Figura C.1(b) Esquema geral do programa desenvolvido nesta pesquisa.

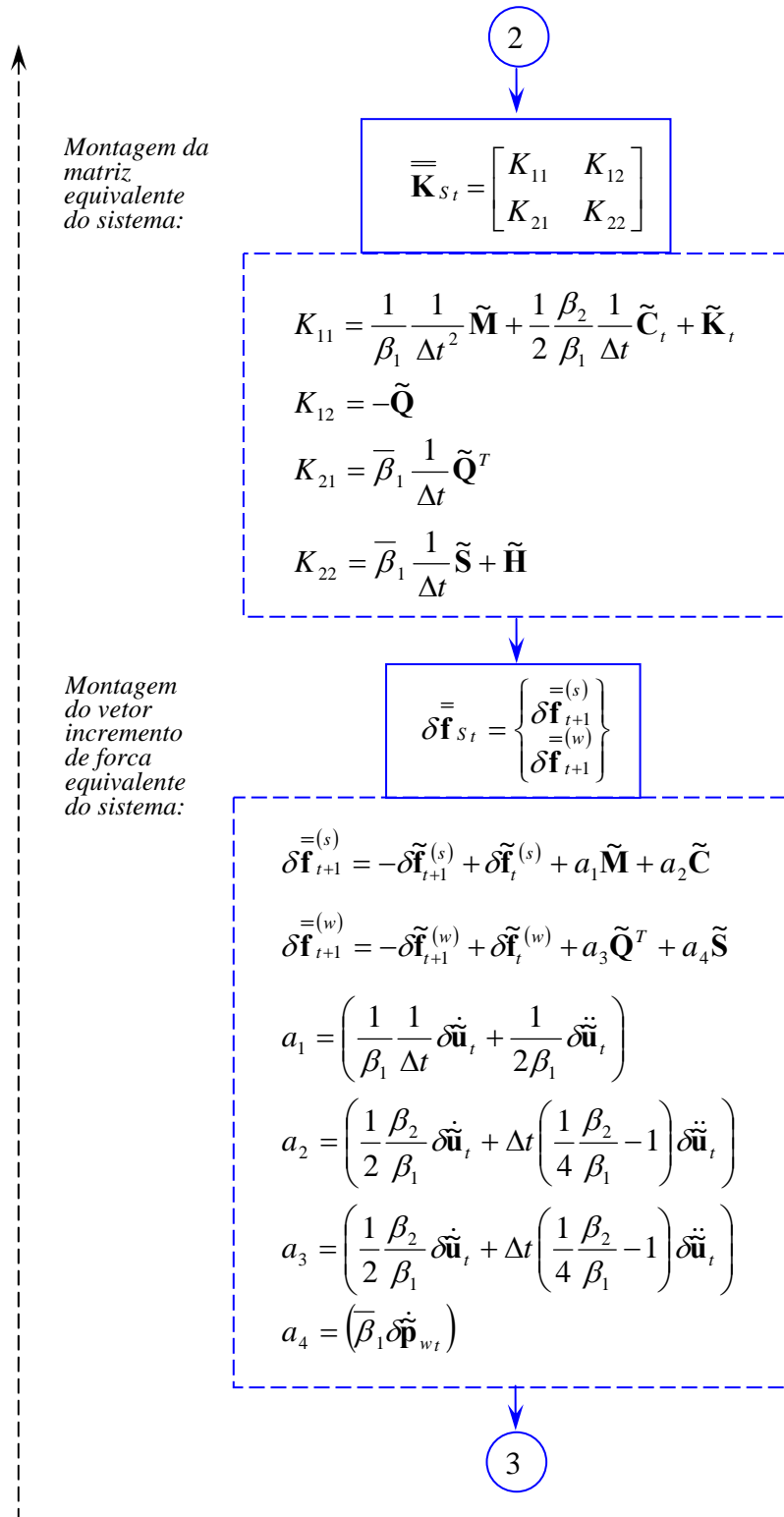


Figura C.1(c) Esquema geral do programa desenvolvido nesta pesquisa.

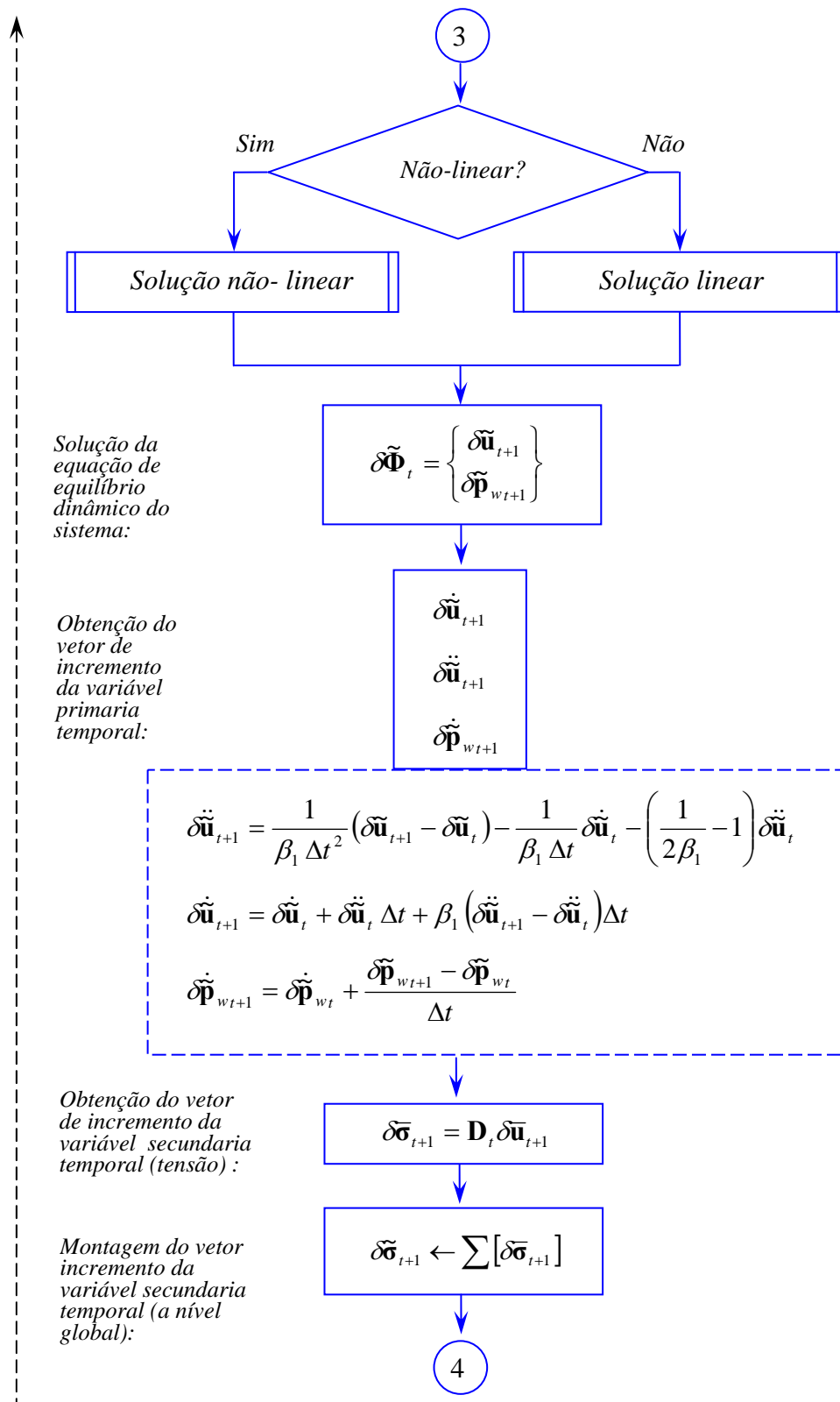


Figura C.1(d) Esquema geral do programa desenvolvido nesta pesquisa.

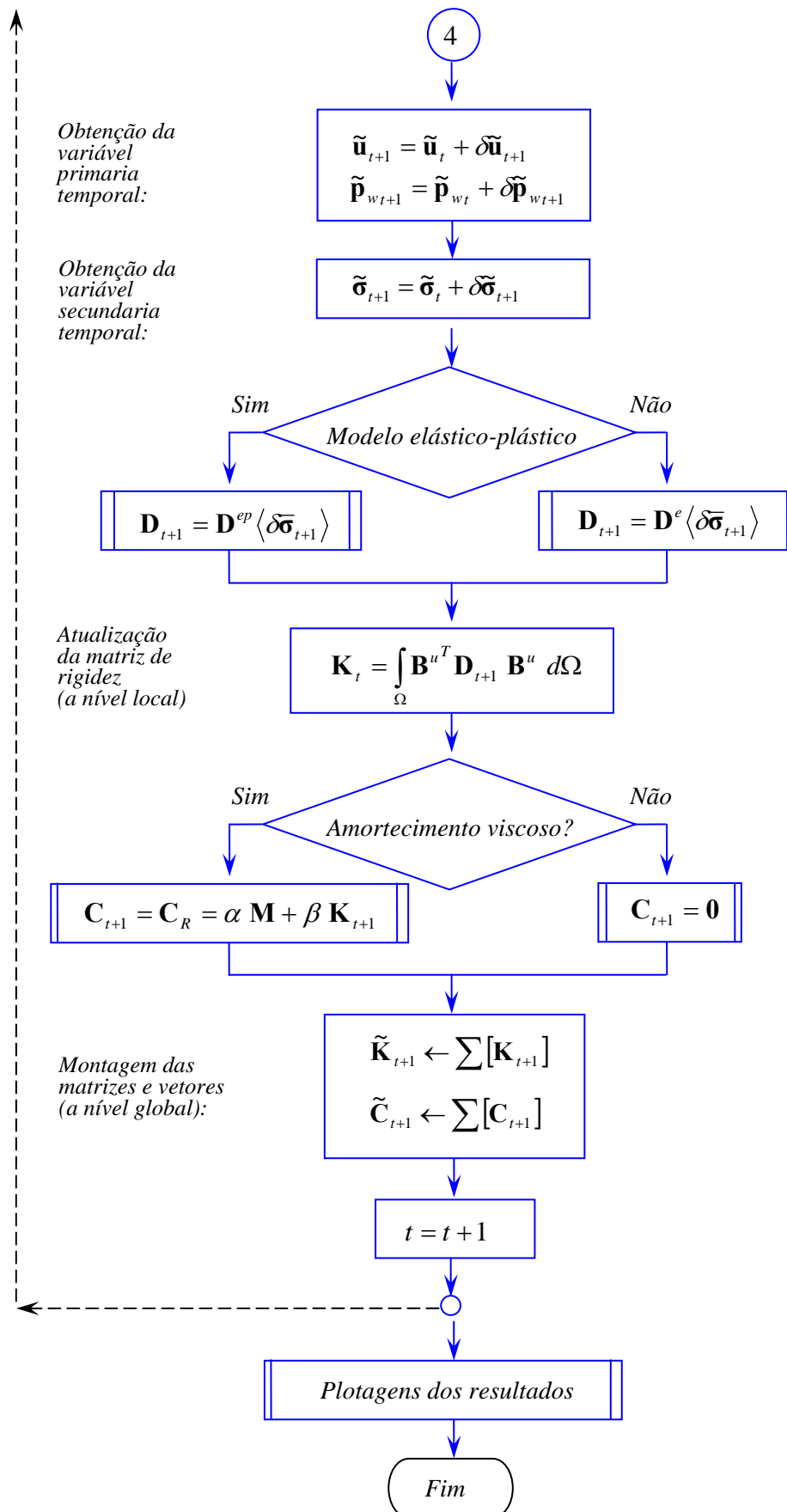


Figura C.1(e) Esquema geral do programa desenvolvido nesta pesquisa.

Apêndices

Apêndice A

Resultados em termos da história dos deslocamentos. Coluna de solo seco. Solução analítica.

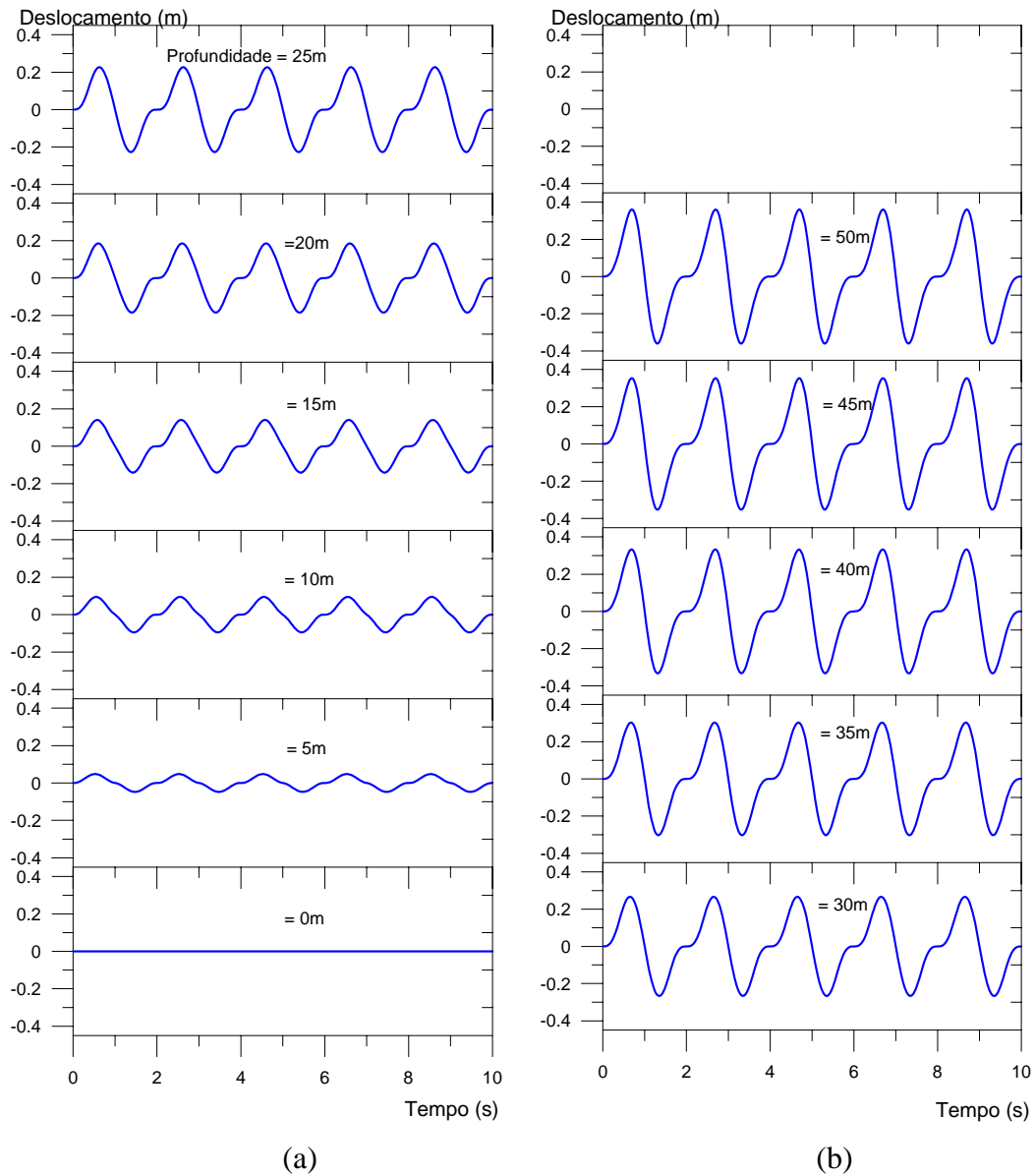


Figura A.1 História dos deslocamentos para a coluna de solo seco. Solução analítica.

Apêndice B
Histórias dos deslocamentos, velocidades e acelerações para a
coluna de solo seco. Solução aproximada [FEM].

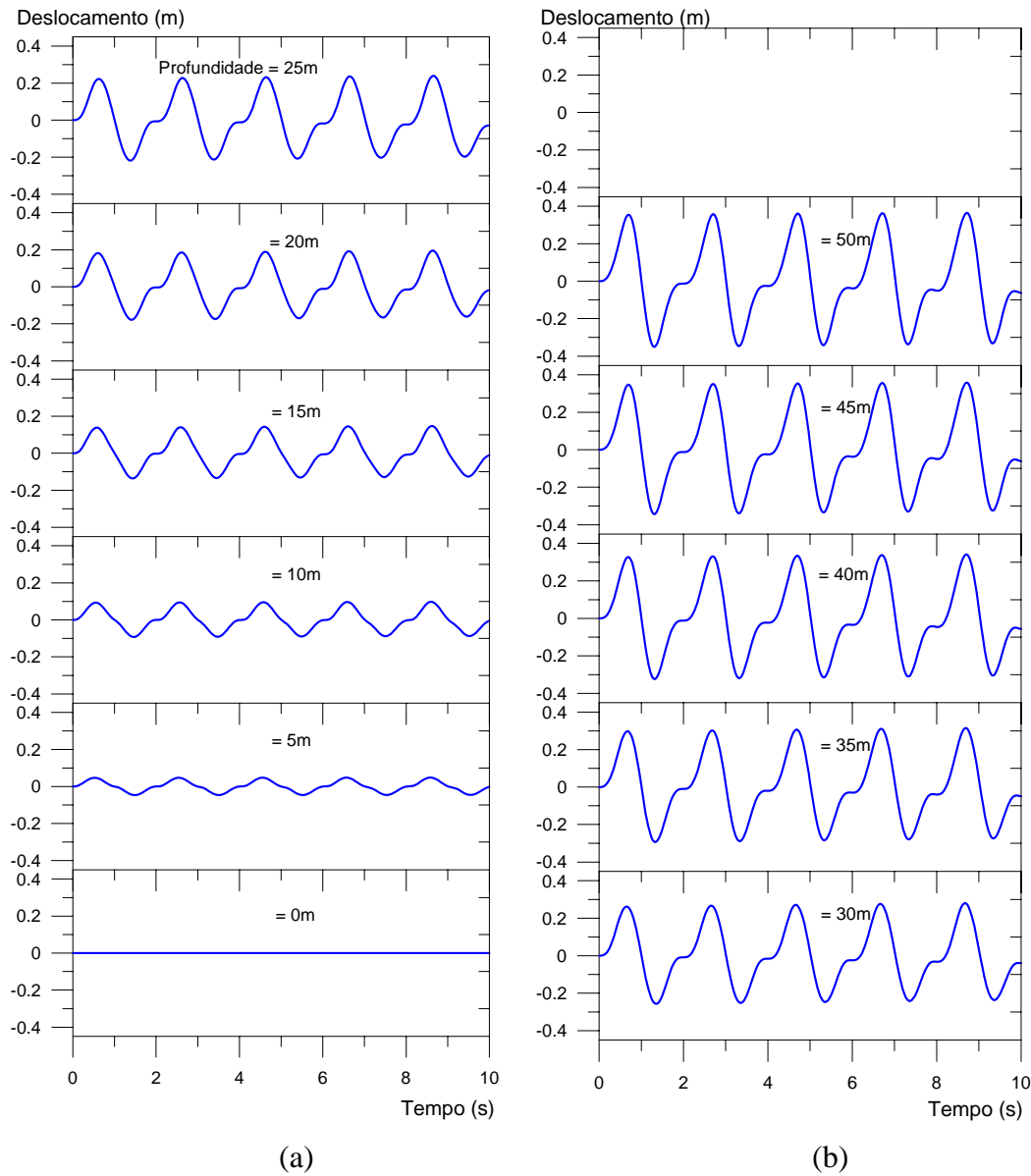
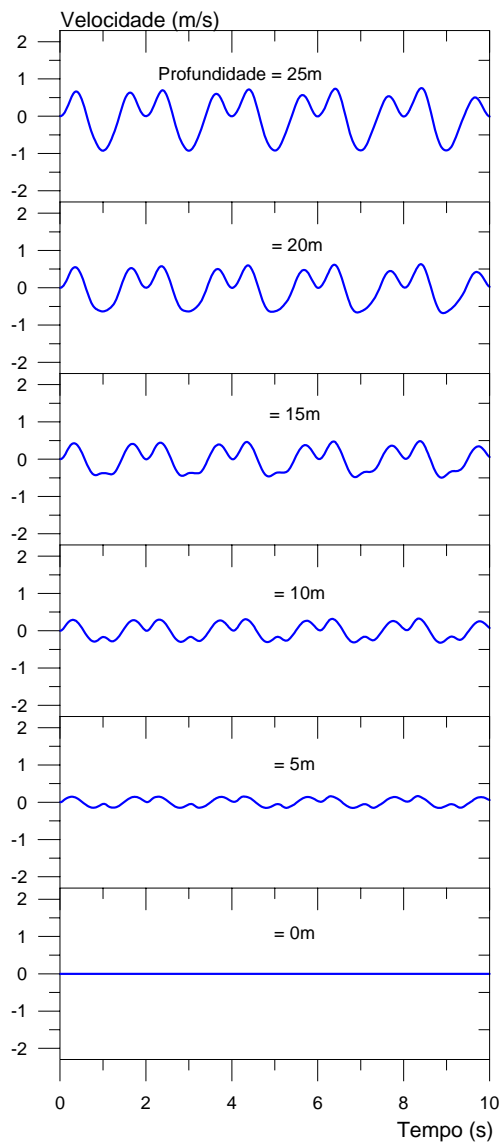
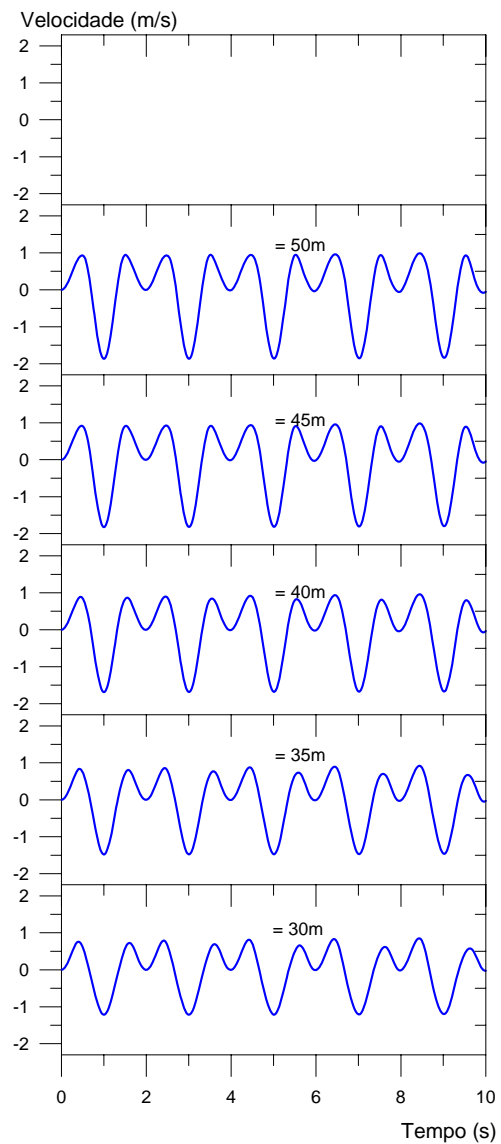


Figura B.1 História dos deslocamentos para coluna de solo seco. Solução aproximada MEF.

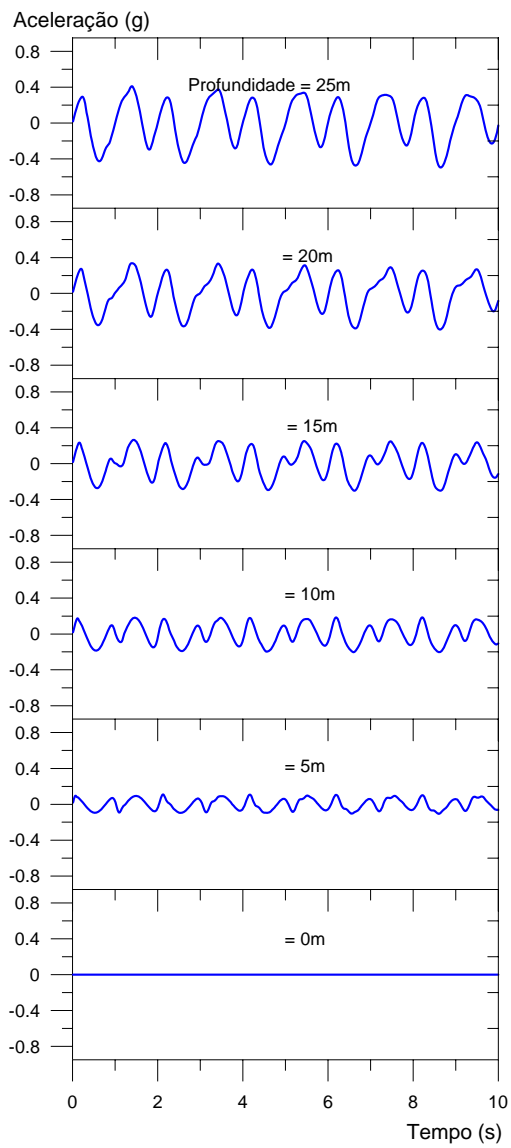


(a)

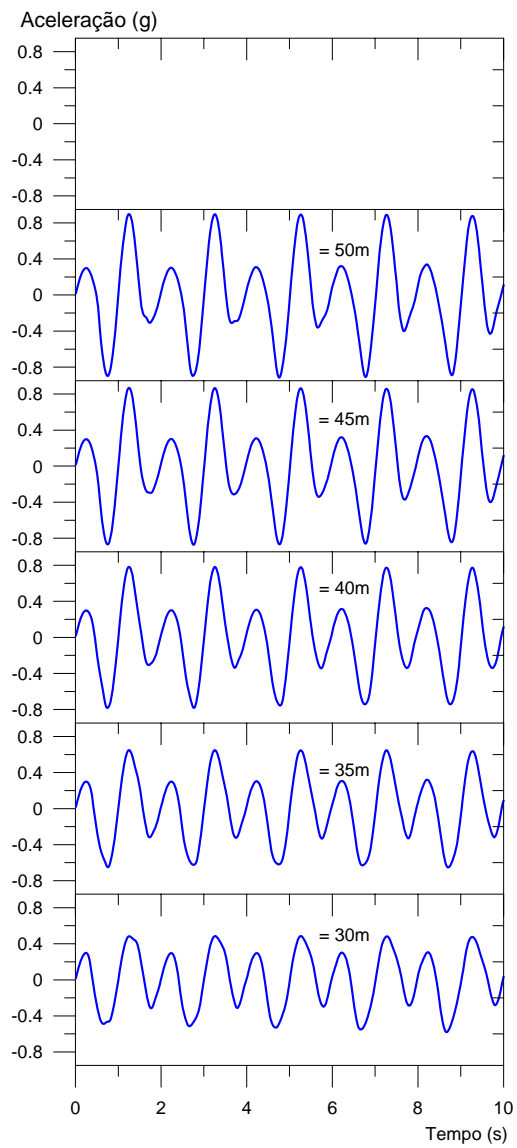


(b)

Figura B.2 História das velocidades para coluna de solo seco. Solução aproximada MEF.



(a)



(b)

Figura B.3 História das acelerações para coluna de solo seco. Solução aproximada MEF.

Apêndice C

Histórias dos deslocamentos, incrementos de poropressão e de acelerações para coluna de solo saturado. Amplitude da onda excitante 0,35g.

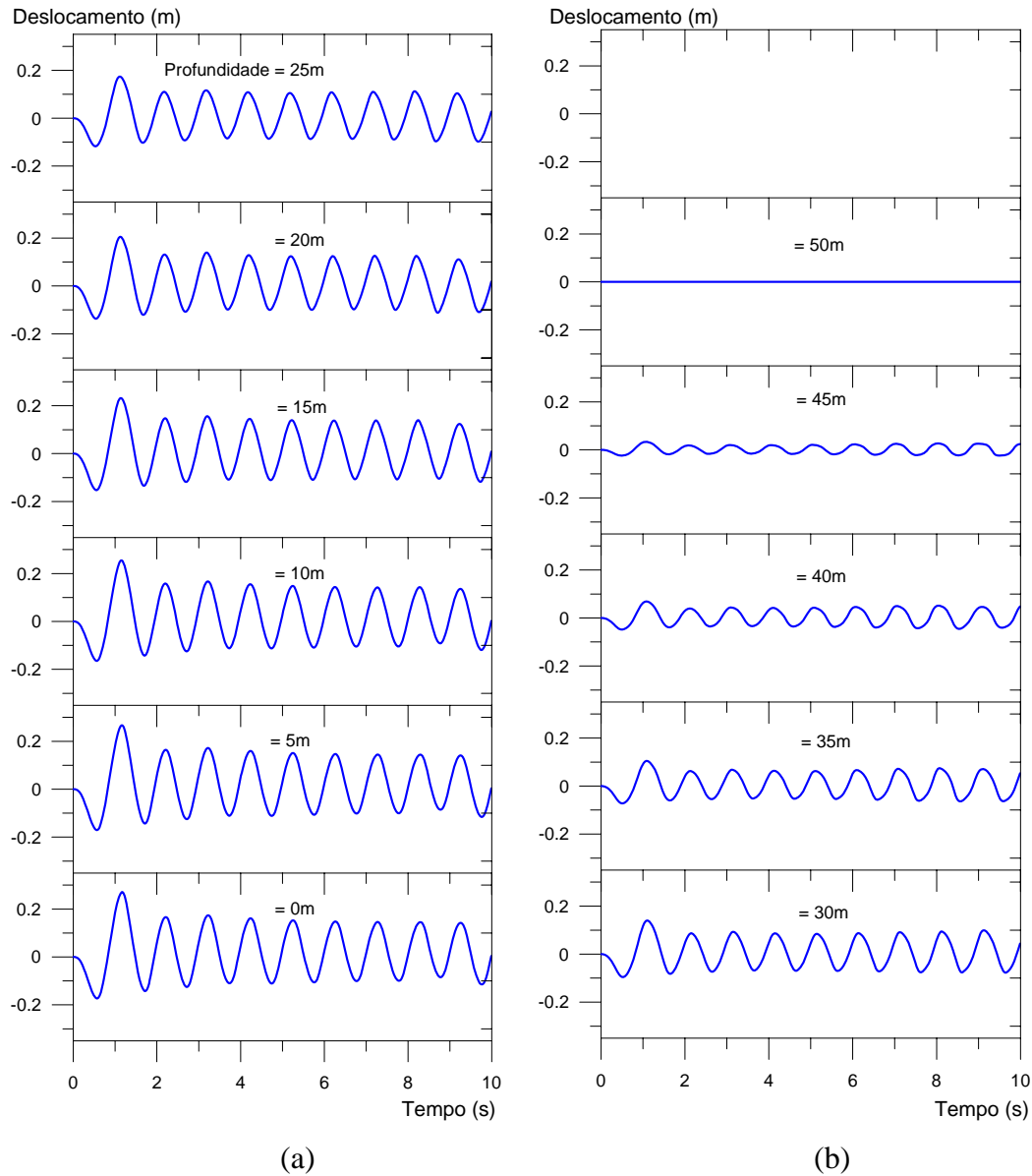


Figura C.1 História dos deslocamentos para coluna de solo saturado. Solução aproximada MEF considerando aceleração horizontal máxima 0,35g .

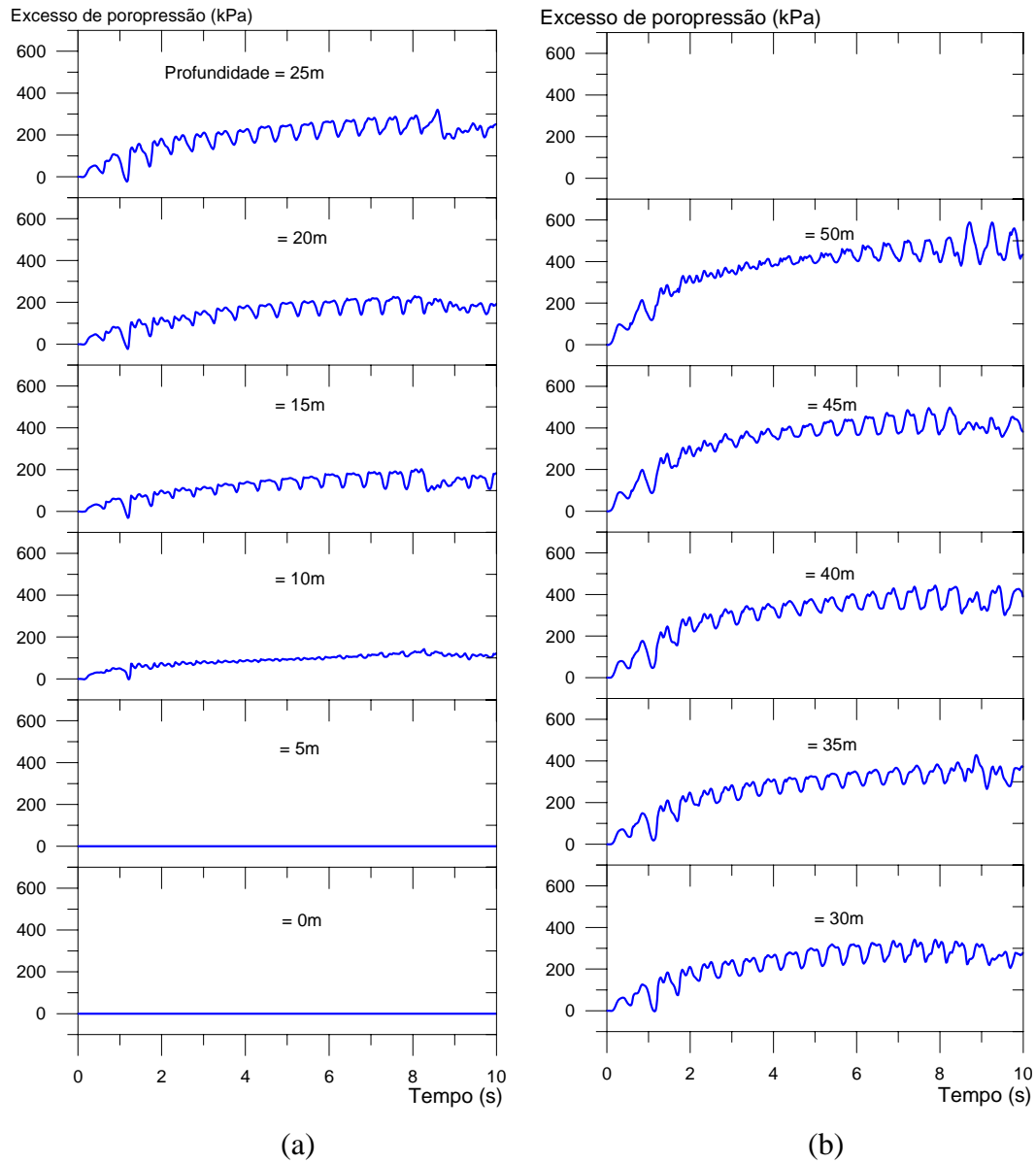


Figura C.2 História dos incrementos de poropressão para coluna de solo saturado. Solução aproximada MEF. Aceleração horizontal máxima 0,35g .

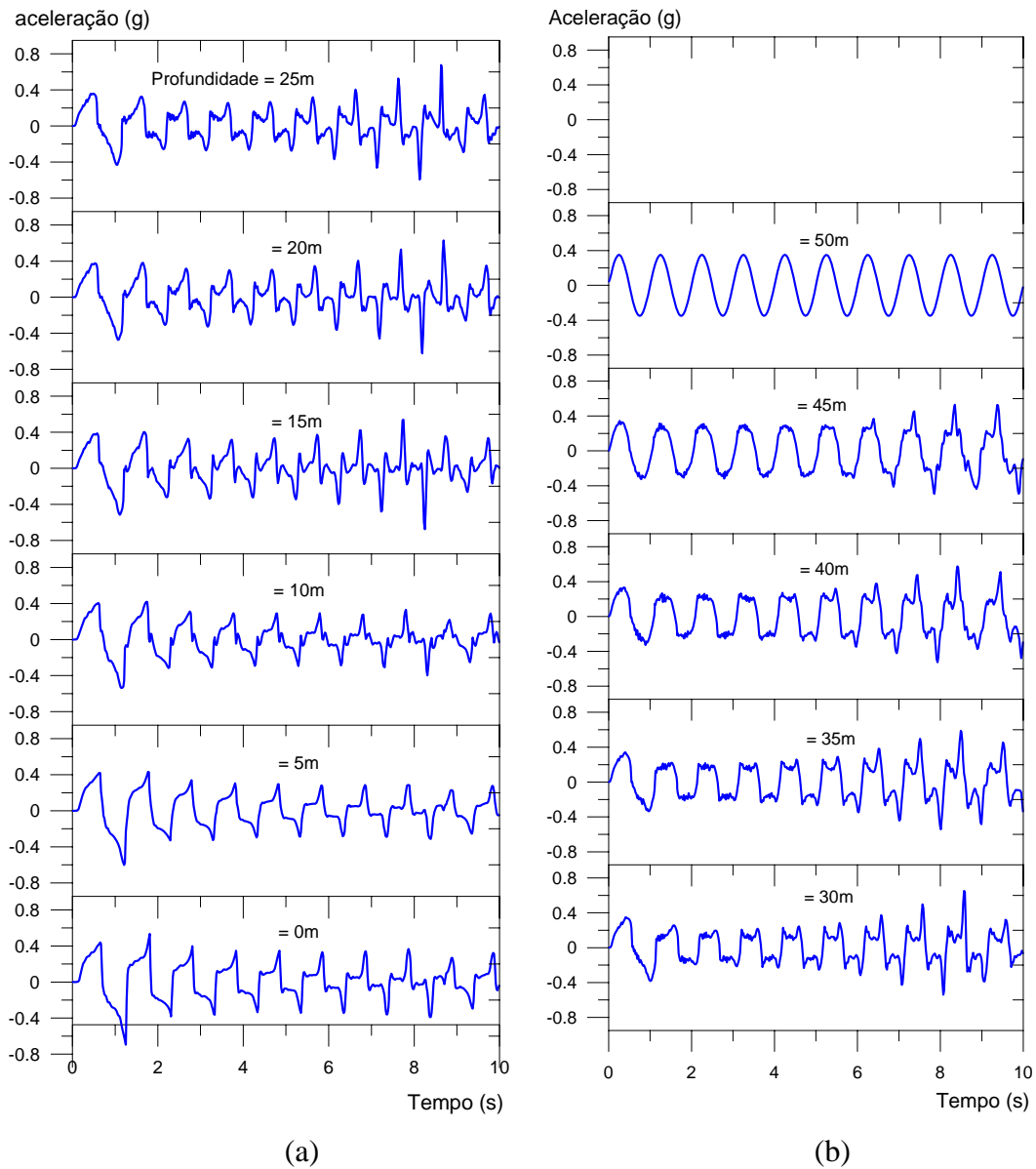


Figura C.3 História das acelerações para coluna de solo saturado. Solução aproximada MEF. Amplitude do registro de aceleração da onda excitante 0,35g .

Apêndice D
Registro dos deslocamentos, incrementos de poropressão e
acelerações para coluna de solo saturado. Aceleração horizontal
máxima 0,40g.

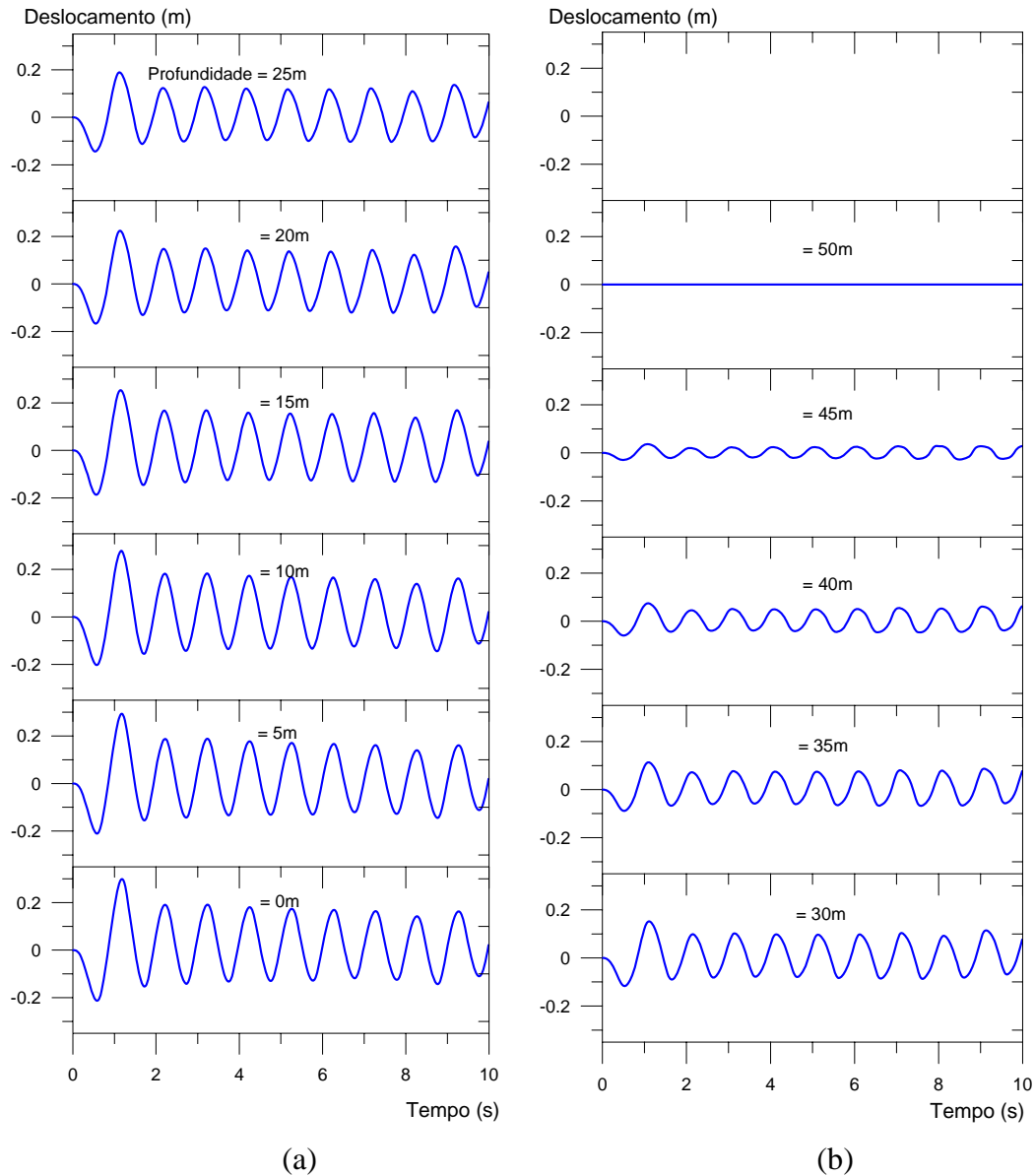


Figura D.1 História dos deslocamentos para coluna de solo saturado. Solução aproximada MEF. Aceleração horizontal máxima 0,40g .

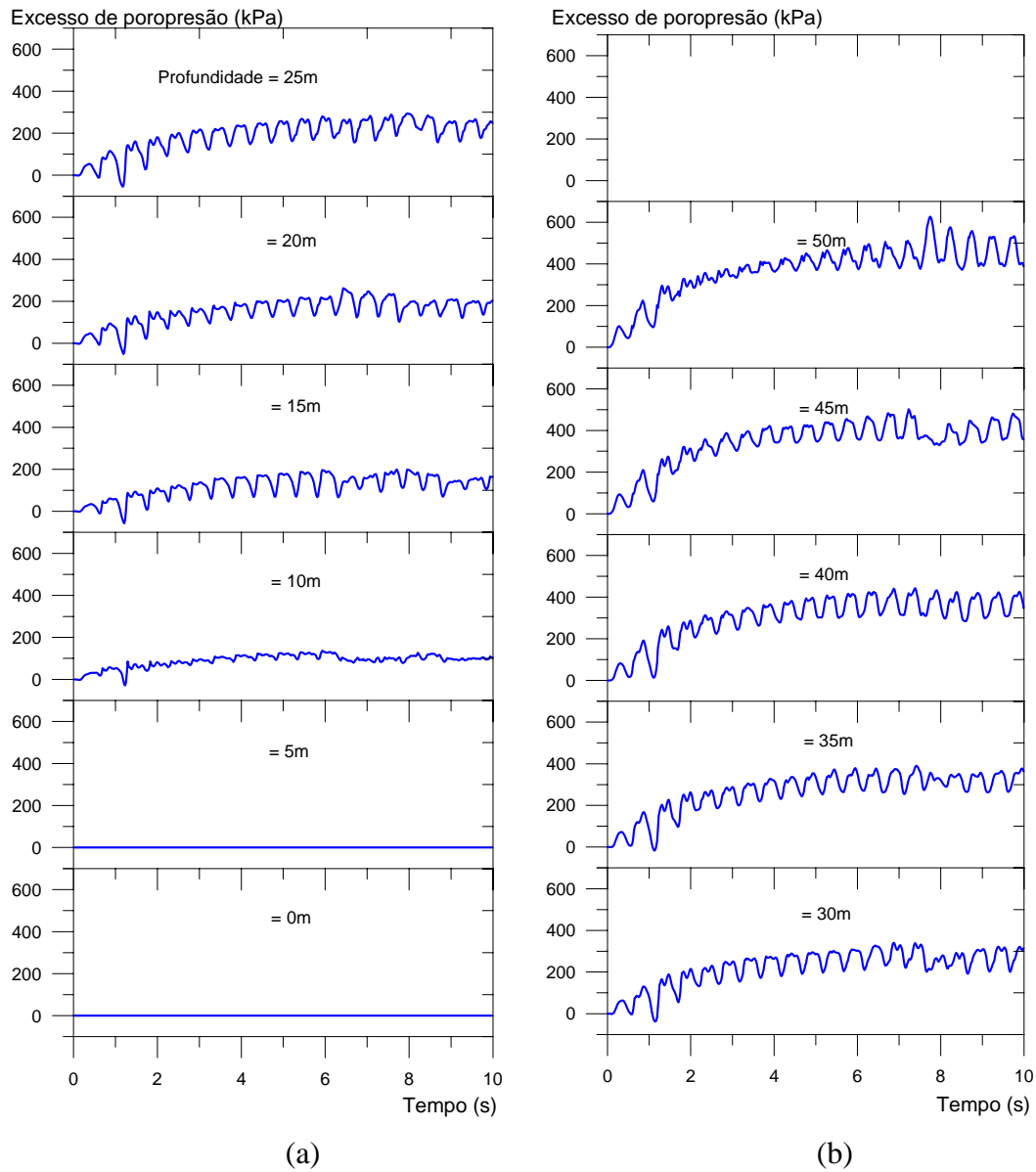


Figura D.2 História dos incremento de poropressão para coluna de solo saturado. Solução aproximada MEF. Aceleração horizontal máxima 0,40g .

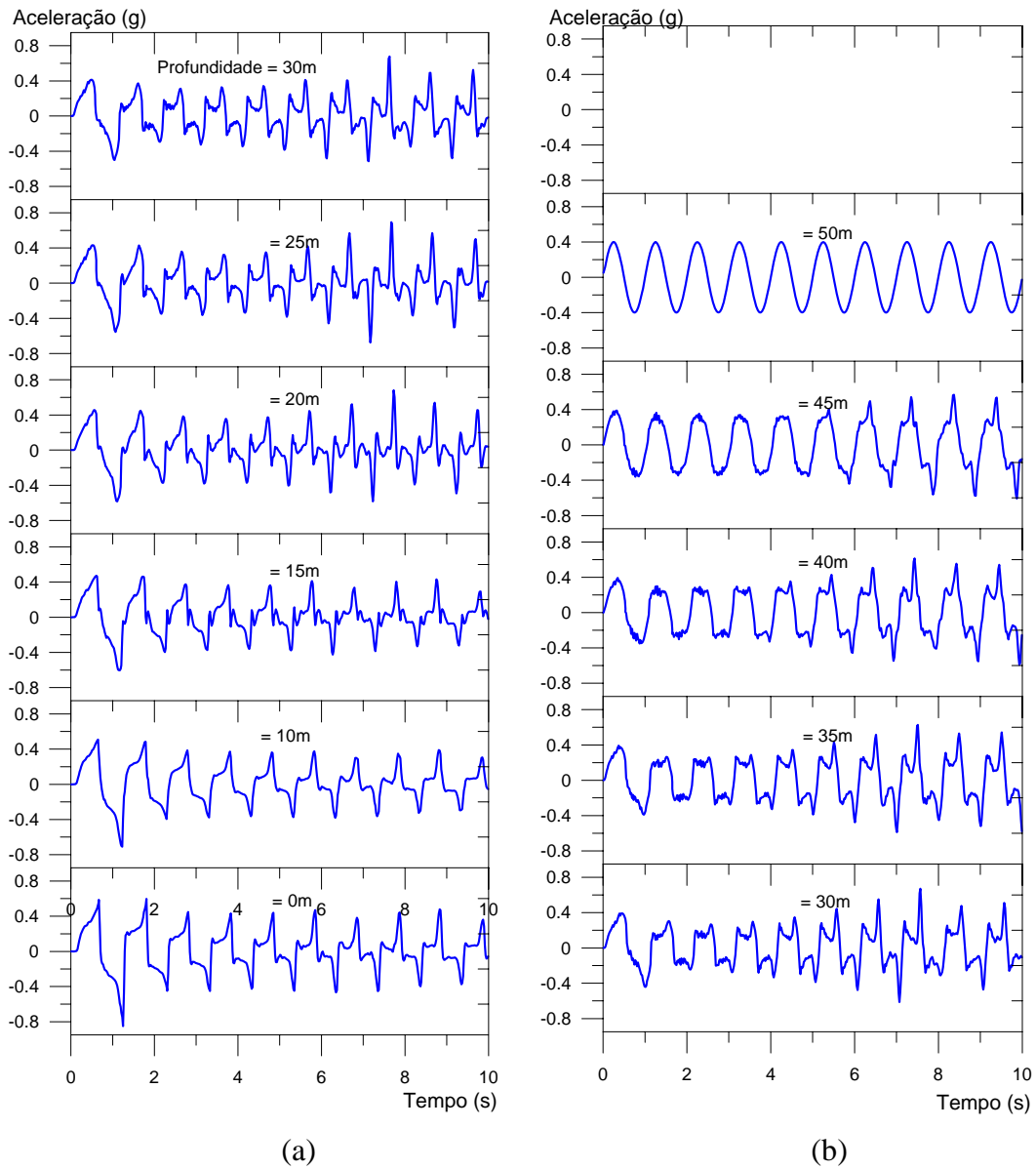


Figura D.3 - História das acelerações para coluna de solo saturado. Solução aproximada MEF. Aceleração horizontal máxima 0,40g .