

Referências Bibliográficas

[1] A.W.Love, *Electromagnetic Horn antennas*, IEEE Press, New York, 1976.

[2] R.E.Collin, *Introduction to Some Fundamental Experiments With Wave Guides*, Proceedings of the IEEE, v87, Issue3, Mar 1999, pp. 508-514.

[3] W.L.Barrow and L.J.Chu, *Theory of Electromagnetic Horns*, Proc. Inst. Radio Eng., 1939, v.27, pp. 51-64.

[4] A.D.Olver, P.J.Clarricoats, A.A.Kishk and L.Shafai, *Microwave Horns and Feeds*, IEE Electromagnetic Wave Series 39, IEE Press, New York, 1994.

[5] P. J. B. Clarricoats and A. D. Olver. *Corrugated Horns for Microwave Antennas*, IEE – Peter Peregrinus Ltd, London – UK, 1984.

[6] P. J. B. Clarricoats, J. R. Descardecı and A. D. Olver, “A novel crosspolar broadband hybrid mode feed”, USNC National Radio Science Meeting – 1988, Colorado, Proceedings of the USNC/URSI National Radio Science Meeting, 1988, v.1, pp. 13-13.

[7] P. J. Clarricoats, J. R. Descardecı and A. D. Olver, “A Low Crosspolar Feed for Broadband Application”, 1988-URSI-International Symp. On Electromagnetic Theory, Thessaloniki USRI, 1988, v.1. pp. 110-119.

[8] A. V. Ghiner and G. I. Surdutovich, “Method of Integral Equations and an Extinction Theorem for Two-Dimensional Problems in Nonlinear Optics”, Physical Review A, vol 50, N.1, July 1994, pp. 714 – 723.

[9] C. A. Balanis. *Advanced Engineering Electromagnetic*. John Wiley & Sons. 1989, USA.

[10] C. A. Balanis. *Antenna Theory. Analysis and Design*. 3rd. ed. John Wiley & Sons. 2005, USA.

[11] A. C. Ludwig. "The Definition of Cross Polarization", IEEE Transactions on Antennas and Propagation, January 1973.

[12] R.F. Harrington. *Time-Harmonic Electromagnetic Fields*. John Wiley & Sons, 2001, USA.

[13] Matrix Operations, Schaum Collection, 1989.

[14] M. Abramovitz., I. A. Stegun. Handbook of mathematical functions. 10^a ed., NBS, 1972.

[15] K. R. Carver, J. W. Mink, "Microstrip Antenna Technology", IEEE Transactions on Antennas and Propagation, vol.29, No. 1, January, 1981.

[16] R. C. Johnson, Antenna Engineering Handbook, 3rd edition, 1993, USA, Mc Graw Hill.

[17] P. D. Potter, "A New horn antenna with suppressed sidelobes and equal beamwidth", Microwave J., pp. 71-78, June 1963.

APENDICE A

Obtenção dos Potenciais Vetores elétrico e magnético de Hertz [9].

A.1

Potencial Vetor de Hertz do Tipo Elétrico

Da equação de Maxwell, $\nabla \cdot \bar{B} = 0$, para um material não magnético, com $\mu = \mu_0$ = permeabilidade do espaço livre, tem-se $\nabla \cdot \bar{B} = \mu_0 \nabla \cdot \bar{H}$. Tendo em vista a identidade vetorial $\nabla \cdot (\nabla \times \bar{\Pi}_e) = 0$, pode-se definir o potencial vetor de Hertz do tipo elétrico a partir de:

$$\bar{H} = j\omega\epsilon_0\bar{\epsilon}_r\nabla\times\bar{\Pi}_e \quad (\text{A.1})$$

Usando-se a equação de Maxwell:

$$\nabla\times\bar{E} = -j\omega\mu_0\bar{H} \quad (\text{A.2})$$

Substituindo-se a equação (A-1) em (A-2), tem-se:

$$\nabla\times\bar{E} = \omega^2\mu_0\epsilon_0\bar{\epsilon}_r\nabla\times\bar{\Pi}_e \quad (\text{A.3})$$

Tendo em vista a identidade vetorial $\nabla\times(\nabla\times\bar{\phi}_e) = 0$, pode-se escrever:

$$\bar{E} = \omega^2\mu_0\epsilon_0\bar{\epsilon}_r\bar{\Pi}_e + \nabla\bar{\phi}_e \quad (\text{A.4})$$

Ficando assim definido o potencial escalar de Hertz do tipo elétrico $\bar{\phi}_e$. Substituindo-se os campos magnético e elétrico, dados por (A.1) e (A.4), respectivamente, na equação de Maxwell:

$$\nabla\times\bar{H} = j\omega\epsilon_0\bar{\epsilon}_r\bar{E} \quad (\text{A.5})$$

obtem-se:

$$(j\omega\epsilon_0\bar{\epsilon}_r)\nabla\times\nabla\times\bar{\Pi}_e = (j\omega\epsilon_0\bar{\epsilon}_r)[\omega^2\mu_0\epsilon_0\bar{\epsilon}_r\bar{\Pi}_e + \nabla\bar{\phi}_e] \quad (\text{A.6})$$

Lembrando que:

$$k_0^2 = \omega^2 \mu_0 \varepsilon_0 \quad (\text{A.7})$$

é o número de onda no espaço livre, pode-se utilizar a identidade vetorial

$$\nabla_x \nabla_x \bar{\Pi}_e = -\nabla^2 \bar{\Pi}_e + \nabla \nabla \cdot \bar{\Pi}_e \quad (\text{A.8})$$

resultando em:

$$-\nabla^2 \bar{\Pi}_e + \nabla \nabla \cdot \bar{\Pi}_e = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e + \nabla \phi_e \quad (\text{A.9})$$

ou

$$\nabla^2 \bar{\Pi}_e + k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e = \nabla \nabla \cdot \bar{\Pi}_e - \nabla \phi_e \quad (\text{A.10})$$

Para que o potencial vetor de Hertz do tipo elétrico, definido por (A.1) seja unívoco, é necessário ainda, uma condição restritiva do tipo do calibre de Lorentz. Impondo-se o calibre de Lorentz:

$$\nabla \nabla \cdot \bar{\Pi}_e - \nabla \phi_e = 0 \quad (\text{A.11})$$

A equação (A.9) se reduz à equação homogênea de Helmholtz,

$$\nabla^2 \bar{\Pi}_e + k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e = 0 \quad (\text{A.12})$$

De (A.1), (A.4) e (A.9) os campos elétrico \bar{E} e magnético \bar{H} , podem ser expressos em termos de $\bar{\Pi}_e$ como seguem:

$$\bar{H} = j\omega \varepsilon_0 \bar{\varepsilon}_r \nabla_x \bar{\Pi}_e \quad (\text{A.13})$$

$$\bar{E} = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e + \nabla \nabla \cdot \bar{\Pi}_e \quad (\text{A.14})$$

A.2

Potencial Vetor de Hertz do Tipo Magnético

Da equação de Maxwell, $\nabla \cdot \bar{D} = 0$, para materiais dielétricos anisotrópicos, com $\bar{D} = \bar{\varepsilon} \bar{E}$, tem-se $\nabla \cdot \bar{\varepsilon} \bar{E} = 0$. Se o material for também homogêneo, então $\nabla \cdot \bar{E} = 0$.

Tendo em vista a identidade vetorial $\nabla \cdot (\nabla_x \bar{\Pi}_h) = 0$, pode-se definir o potencial vetor de Hertz do tipo magnético, a partir de:

$$\bar{E} = -j\omega \mu_0 \nabla_x \bar{\Pi}_h \quad (\text{A.15})$$

Com a equação de Maxwell: $\nabla \times \bar{H} = j\omega \varepsilon_0 \bar{\varepsilon}_r \bar{E}$, obtém-se:

$$\nabla \times \bar{H} = \omega^2 \mu_0 \varepsilon_0 \bar{\varepsilon}_r \nabla \times \bar{\Pi}_h \quad (\text{A.16})$$

Considerando a identidade vetorial $\nabla \times (\nabla \phi_h) = 0$, pode-se escrever de (A.16):

$$\bar{H} = \omega^2 \mu_0 \varepsilon_0 \bar{\varepsilon}_r \bar{\Pi}_h + \nabla \phi_h \quad (\text{A.17})$$

Ficando assim definido o potencial escalar de Hertz do tipo magnético ϕ_h .

Substituindo \bar{E} e \bar{H} , dados por (A.15) e (A.17), respectivamente, na equação de Maxwell: $\nabla \times \bar{E} = -j\omega \mu_0 \bar{H}$, obtém-se:

$$(j\omega \mu_0) \nabla \times \nabla \times \bar{\Pi}_h = (j\omega \mu_0) [\omega^2 \mu_0 \varepsilon_0 \bar{\varepsilon}_r \bar{\Pi}_h + \nabla \phi_h] \quad (\text{A.18})$$

com $k_0^2 = \omega^2 \mu_0 \varepsilon_0$, então:

$$\nabla \times \nabla \times \bar{\Pi}_h = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_h + \nabla \phi_h \quad (\text{A.19})$$

Utilizando a identidade vetorial $\nabla \times \nabla \times \bar{\Pi}_h = -\nabla^2 \bar{\Pi}_h + \nabla \nabla \cdot \bar{\Pi}_h$ em (A.19), resulta em:

$$\nabla^2 \bar{\Pi}_h + k_0^2 \bar{\varepsilon}_r \bar{\Pi}_h = \nabla \nabla \cdot \bar{\Pi}_h - \nabla \phi_h \quad (\text{A.20})$$

Para que $\bar{\Pi}_h$, definido por (A.20) seja unívoco, impõe-se uma condição restritiva do tipo calibre de Lorentz:

$$\nabla \nabla \cdot \bar{\Pi}_h - \nabla \phi_h = 0 \quad (\text{A.21})$$

resultando:

$$\nabla^2 \bar{\Pi}_h + k_0^2 \bar{\varepsilon}_r \bar{\Pi}_h = 0 \quad (\text{A.22})$$

De (A.15), (A.17), (A.21) e (A.22) os campos podem ser expressos em termos de $\bar{\Pi}_h$, como segue:

$$\bar{E} = -j\omega \mu_0 \nabla \times \bar{\Pi}_h \quad (\text{A.23})$$

$$\bar{H} = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_h + \nabla \nabla \cdot \bar{\Pi}_h = \nabla \nabla \cdot \bar{\Pi}_h - \nabla^2 \bar{\Pi}_h \quad (\text{A.24})$$

A.3

Solução Geral com os Potenciais Vetores de Hertz dos Tipos Elétrico e Magnético

Para a solução de um problema usando os potenciais vetores de Hertz do tipo elétrico e magnético, $\bar{\Pi}_e$ e $\bar{\Pi}_h$, respectivamente, a primeira providência é obter as soluções das equações (A.12) e (A.22). Conhecidos $\bar{\Pi}_e$ e $\bar{\Pi}_h$, os campos elétrico e magnético são obtidos como as somas de (A.14) e (A.23) para o campo elétrico; e (A.13) e (A.24) para o campo magnético:

$$\bar{E} = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e + \nabla \nabla \cdot \bar{\Pi}_e - j\omega \mu_0 \nabla \times \bar{\Pi}_h \quad (\text{A.25})$$

$$\bar{H} = \nabla \nabla \cdot \bar{\Pi}_h - \nabla^2 \bar{\Pi}_h + j\omega \varepsilon_0 \bar{\varepsilon}_r \nabla \times \bar{\Pi}_e \quad (\text{A.26})$$

com

$$\bar{\varepsilon}_r = \begin{pmatrix} \varepsilon_t & 0 & 0 \\ 0 & \varepsilon_t & 0 \\ 0 & 0 & \varepsilon_z \end{pmatrix}$$

onde:

ε_t é a componente tangencial do tensor da permissividade relativa do meio,

ε_z é a componente axial do tensor da permissividade relativa do meio.

APÊNDICE B

Determinação dos Campos Elétrico e Magnético a partir dos Potenciais Vetores de Hertz

Para o potencial vetor magnético, tem-se que [9]:

$$\nabla^2 \bar{\Pi}_h + k_0^2 \bar{\epsilon}_r \bar{\Pi}_h = 0 \quad (\text{B.1})$$

ou

$$\nabla^2 \bar{\Pi}_h + k_0^2 \begin{pmatrix} \epsilon_r & 0 & 0 \\ 0 & \epsilon_r & 0 \\ 0 & 0 & \epsilon_z \end{pmatrix} \bar{\Pi}_h = 0 \quad (\text{B.2})$$

Como $\bar{\Pi}_h$ só tem componente na direção z, então:

$$\nabla^2 \bar{\Pi}_h = \hat{z} \left\{ \frac{\partial^2 \bar{\Pi}_h}{\partial r^2} + \frac{1}{r} \frac{\partial \bar{\Pi}_h}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \bar{\Pi}_h}{\partial \phi^2} + \frac{\partial^2 \bar{\Pi}_h}{\partial z^2} \right\} \quad (\text{B.3})$$

logo

$$\frac{\partial^2 \Pi_h}{\partial r^2} + \frac{1}{r} \frac{\partial \Pi_h}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Pi_h}{\partial \phi^2} + \frac{\partial^2 \Pi_h}{\partial z^2} + k_0^2 \epsilon_z \Pi_h = 0 \quad (\text{B.4})$$

Assumindo uma propagação na forma $e^{-\gamma z} \Rightarrow$

$$\frac{\partial^2 \Pi_h}{\partial z^2} = \gamma_z^2 \Pi_h,$$

obtem-se:

$$\frac{\partial^2 \Pi_h}{\partial r^2} + \frac{1}{r} \frac{\partial \Pi_h}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Pi_h}{\partial \phi^2} + k_{CH}^2 \Pi_h = 0 \quad (\text{B.5})$$

chamando

$$k_{CH}^2 = k_0^2 \epsilon_z + \gamma_z^2 \quad (\text{B.6})$$

Então, resolvendo a equação de onda por substituição de variáveis, chega-se a:

$$\bar{\Pi}_h = \hat{z} \begin{cases} B_n J_n(k_{CH} r) e^{-\gamma z} \text{sen}(n\phi), 0 \leq r \leq r_1 \\ [E_n J_n(K_1 r) + F_n Y_n(K_1 r)] e^{-\gamma z} \text{sen}(n\phi), r_1 \leq r \leq r_0 \end{cases} \quad (\text{B.7})$$

Analisando o potencial vetor elétrico [9], tem-se:

$$\nabla^2 \bar{\Pi}_e + k_0^2 \bar{\varepsilon}_r \bar{\Pi}_e = 0 \quad (\text{B.8})$$

Como $\bar{\Pi}_e$ só tem componente na direção z, então:

$$\bar{\Pi}_e = \Pi_e \hat{z} = \begin{bmatrix} 0\hat{r} \\ 0\hat{\phi} \\ \Pi_e \hat{z} \end{bmatrix} \quad (\text{B.9})$$

$$\nabla^2 \bar{\Pi}_e = \hat{z} \left\{ \frac{\partial^2 \bar{\Pi}_e}{\partial r^2} + \frac{1}{r} \frac{\partial \bar{\Pi}_e}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \bar{\Pi}_e}{\partial \phi^2} + \frac{\partial^2 \bar{\Pi}_e}{\partial z^2} \right\} \quad (\text{B.10})$$

Assumindo uma propagação na forma $e^{-\gamma z}$, tem-se:

$$\frac{\partial^2 \Pi_e}{\partial z^2} = \gamma_z^2 \Pi_e \quad (\text{B.11})$$

então:

$$\frac{\partial^2 \Pi_e}{\partial r^2} + \frac{1}{r} \frac{\partial \Pi_e}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Pi_e}{\partial \phi^2} + k_{CE}^2 \Pi_e = 0 \quad (\text{B.12})$$

com

$$k_{CE}^2 = k_0^2 \varepsilon_z + \gamma_z^2 \quad (\text{B.13})$$

Portanto:

$$\bar{\Pi}_e = \hat{z} \{ A_n J_n(k_{CE} r) e^{-\gamma z} \cos(n\phi) \}, 0 \leq r \leq r_1 \quad (\text{B.14})$$

$$\bar{\Pi}_e = \hat{z} \{ (C_n J_n(K_1 r) + D_n Y_n(K_1 r)) e^{-\gamma z} \cos(n\phi) \}, r_1 \leq r \leq r_0 \quad (\text{B.15})$$

O campo elétrico, obtido através dos potenciais vetores é dado por [9]:

$$\vec{E} = k_0^2 \varepsilon_z \bar{\Pi}_e + \nabla(\nabla \cdot \bar{\Pi}_e) - j\omega\mu_0 \nabla \times \bar{\Pi}_h \quad (\text{B.16})$$

$$\nabla \cdot \bar{\Pi}_e = \nabla \cdot (\Pi_e \hat{z}) = \frac{\partial \Pi_e}{\partial z} \quad (\text{B.17})$$

$$\nabla(\nabla \cdot \bar{\Pi}_e) = \hat{r} \frac{\partial^2 \Pi_e}{\partial r \partial z} + \hat{\phi} \frac{1}{r} \frac{\partial^2 \Pi_e}{\partial \phi \partial z} + \hat{z} \frac{\partial^2 \Pi_e}{\partial z^2} \quad (\text{B.18})$$

$$\nabla \times \bar{\Pi}_h = \hat{r} \frac{1}{r} \frac{\partial \Pi_h}{\partial \phi} - \hat{\phi} \frac{\partial \Pi_h}{\partial r} \quad (\text{B.19})$$

Desta forma, substituindo (B.18) e (B.19) em (B.16) pode-se reescrever o campo elétrico como:

$$\vec{E} = k_0^2 \varepsilon_z \bar{\Pi}_e \hat{z} + \left\{ \hat{r} \frac{\partial^2 \Pi_e}{\partial r \partial z} + \hat{\phi} \frac{1}{r} \frac{\partial^2 \Pi_e}{\partial \phi \partial z} + \hat{z} \frac{\partial^2 \Pi_e}{\partial z^2} \right\} - j\omega\mu_0 \left(\hat{r} \frac{1}{r} \frac{\partial \Pi_h}{\partial \phi} - \hat{\phi} \frac{\partial \Pi_h}{\partial r} \right) \quad (\text{B.20})$$

$$\vec{E} = E_r \hat{r} + E_\phi \hat{\phi} + E_z \hat{z} \quad (\text{B.21})$$

com

$$E_r = \frac{\partial^2 \Pi_e}{\partial r \partial z} - j\omega\mu \frac{1}{r} \frac{\partial \Pi_h}{\partial \phi} \quad (\text{B.22})$$

$$E_\phi = \frac{1}{r} \frac{\partial^2 \Pi_e}{\partial \phi \partial z} + j\omega\mu \frac{\partial \Pi_h}{\partial r} \quad (\text{B.23})$$

$$E_z = \frac{\partial^2 \Pi_e}{\partial z^2} + k_0^2 \varepsilon_z \Pi_e \quad (\text{B.24})$$

Analisando a equação (B.24) junto com a equação (B.13), verifica-se que:

$$\begin{cases} E_z^i = [k_0^2 \varepsilon_z + \gamma_z^2] \Pi_e = k_{CE}^2 \Pi_e, 0 \leq r \leq r_1 \\ E_z^o = [k_0^2 + \gamma_z^2] \Pi_e = K_1^2 \Pi_e, r_1 \leq r \leq r_0 \end{cases} \quad (\text{B.25})$$

onde $\gamma_z^2 = \frac{\partial^2}{\partial z^2}$

então

$$\begin{cases} \Pi_e = \frac{1}{k_{CE}^2} E_z^i, 0 \leq r \leq r_1 \\ \Pi_e = \frac{1}{K_1^2} E_z^o, r_1 \leq r \leq r_2 \end{cases} \quad (\text{B.26})$$

O campo magnético, obtido pelos potenciais vetores é dado por:

$$\vec{H} = k_0^2 \bar{\varepsilon}_r \bar{\Pi}_h + \nabla(\nabla \cdot \bar{\Pi}_h) + j\omega\varepsilon_0 \bar{\varepsilon}_r \nabla \times \bar{\Pi}_e \quad (\text{B.27})$$

$$\nabla \cdot \bar{\Pi}_h = \nabla \cdot (\bar{\Pi}_h \hat{z}) = \frac{\partial \Pi_h}{\partial z} \quad (\text{B.28})$$

$$\nabla(\nabla \cdot \bar{\Pi}_h) = \hat{r} \frac{\partial^2 \Pi_h}{\partial r \partial z} + \hat{\phi} \frac{1}{r} \frac{\partial^2 \Pi_h}{\partial \phi \partial z} + \hat{z} \frac{\partial^2 \Pi_h}{\partial z^2} \quad (\text{B.29})$$

$$\nabla \times \bar{\Pi}_e = \hat{r} \frac{1}{r} \frac{\partial \Pi_e}{\partial \phi} - \hat{\phi} \frac{\partial \Pi_e}{\partial r} \quad (\text{B.30})$$

e

$$\bar{\varepsilon}_r = \begin{pmatrix} \varepsilon_t & 0 & 0 \\ 0 & \varepsilon_t & 0 \\ 0 & 0 & \varepsilon_z \end{pmatrix} \quad (\text{B.31})$$

resultando em:

$$\vec{H} = k_0^2 \varepsilon_z \Pi_h \hat{z} + \left\{ \hat{r} \frac{\partial^2 \Pi_h}{\partial r \partial z} + \hat{\phi} \frac{1}{r} \frac{\partial^2 \Pi_h}{\partial \phi \partial z} + \hat{z} \frac{\partial^2 \Pi_h}{\partial z^2} \right\} + j\omega \varepsilon_0 \varepsilon_t \left(\hat{r} \frac{1}{r} \frac{\partial \Pi_e}{\partial \phi} - \hat{\phi} \frac{\partial \Pi_e}{\partial r} \right) \quad (\text{B.32})$$

$$\vec{H} = H_r \hat{r} + H_\phi \hat{\phi} + H_z \hat{z} \quad (\text{B.33})$$

com

$$H_r = \frac{\partial^2 \Pi_h}{\partial r \partial z} + j\omega \varepsilon_0 \varepsilon_t \frac{1}{r} \frac{\partial \Pi_e}{\partial \phi} \quad (\text{B.34})$$

$$H_\phi = \frac{1}{r} \frac{\partial^2 \Pi_h}{\partial \phi \partial z} - j\omega \varepsilon_0 \varepsilon_t \frac{\partial \Pi_e}{\partial r} \quad (\text{B.35})$$

$$H_z = \frac{\partial^2 \Pi_h}{\partial z^2} + k_0^2 \varepsilon_z \Pi_h \quad (\text{B.36})$$

Analisando as equações (B.36) e (B.6), verifica-se que:

$$\begin{cases} H_z = [\varepsilon_z k_0^2 + \gamma_z^2] \Pi_h = k_{CH}^2 \Pi_h, 0 \leq r \leq r_1 \\ H_z = [k_0^2 + \gamma_z^2] \Pi_h = K_1^2 \Pi_h, r_1 \leq r \leq r_0 \end{cases} \quad (\text{B.37})$$

então:

$$\begin{cases} \Pi_h = \frac{1}{k_{CH}^2} H_z^i, 0 \leq r \leq r_1 \\ \Pi_h = \frac{1}{K_1^2} H_z^o, r_1 \leq r \leq r_0 \end{cases} \quad (\text{B.38})$$

com as equações (B.26) e (B.38) pode-se reescrever as equações (B.22), (B.23), (B.34) e (B.35), para as duas regiões de estudo:

1) região i ($0 \leq r \leq r_1$):

$$E_r^i = \frac{1}{k_{CE}^2} \frac{\partial^2 E_z^i}{\partial r \partial z} - \frac{j\omega \mu_0}{rk_{CE}^2} \frac{\partial H_z^i}{\partial \phi} = \frac{1}{k_{CE}^2} \left\{ \frac{\partial^2 E_z^i}{\partial r \partial z} - \frac{j\omega \mu_0}{r} \frac{\partial H_z^i}{\partial \phi} \right\} \quad (\text{B.39})$$

$$E_\phi^i = \frac{1}{k_{CE}^2} \frac{1}{r} \frac{\partial^2 E_z^i}{\partial \phi \partial z} + \frac{j\omega \mu_0}{k_{CE}^2} \frac{\partial H_z^i}{\partial r} = \frac{1}{k_{CE}^2} \left\{ \frac{1}{r} \frac{\partial^2 E_z^i}{\partial \phi \partial z} + j\omega \mu_0 \frac{\partial H_z^i}{\partial r} \right\} \quad (\text{B.40})$$

$$H_r^i = \frac{1}{k_{CH}^2} \frac{\partial^2 H_z^i}{\partial r \partial z} + \frac{j\omega \varepsilon_0 \varepsilon_{ii}}{k_{CH}^2} \frac{1}{r} \frac{\partial E_z^i}{\partial \phi} = \frac{1}{k_{CH}^2} \left\{ \frac{\partial^2 H_z^i}{\partial r \partial z} + \frac{j\omega \varepsilon_0 \varepsilon_{ii}}{r} \frac{\partial E_z^i}{\partial \phi} \right\} \quad (\text{B.41})$$

$$H_\phi^i = \frac{1}{k_{CH}^2 r} \frac{\partial^2 H_z^i}{\partial \phi \partial z} - \frac{j\omega \varepsilon_0 \varepsilon_{ii}}{k_{CH}^2} \frac{\partial E_z^i}{\partial r} = \frac{1}{k_{CH}^2} \left\{ \frac{1}{r} \frac{\partial^2 H_z^i}{\partial \phi \partial z} - j\omega \varepsilon_0 \varepsilon_{ii} \frac{\partial E_z^i}{\partial r} \right\} \quad (\text{B.42})$$

onde, considerando $\gamma = j\beta$ (desconsiderando a atenuação, $\alpha = 0$):

$$k_{CE}^2 = k_{CH}^2 = K^2 = k_0^2 \varepsilon_{zi} + \gamma_z^2 = k_0^2 \varepsilon_{zi} - \beta_z^2 \quad (\text{B.43})$$

2) região o ($r_1 \leq r \leq r_0$): $\varepsilon_{to} = \varepsilon_{ro} = 1$

$$E_r^o = \frac{1}{K_1^2} \left\{ \frac{\partial^2 E_z^o}{\partial r \partial z} - \frac{j\omega \mu_0}{r} \frac{\partial H_z^o}{\partial \phi} \right\} \quad (\text{B.44})$$

$$E_\phi^o = \frac{1}{K_1^2} \left\{ \frac{1}{r} \frac{\partial^2 E_z^o}{\partial \phi \partial z} + j\omega \mu_0 \frac{\partial H_z^o}{\partial r} \right\} \quad (\text{B.45})$$

$$H_r^o = \frac{1}{K_1^2} \left\{ \frac{\partial^2 H_z^o}{\partial r \partial z} + \frac{j\omega \varepsilon_0 \varepsilon_{to}}{r} \frac{\partial E_z^o}{\partial \phi} \right\} \quad (\text{B.46})$$

$$H_\phi^o = \frac{1}{k_1^2} \left\{ \frac{1}{r} \frac{\partial^2 H_z^o}{\partial \phi \partial z} - j\omega \varepsilon_0 \varepsilon_{to} \frac{\partial E_z^o}{\partial r} \right\} \quad (\text{B.47})$$

e

$$K_1^2 = k_0^2 + \gamma_z^2 = k_0^2 - \beta_z^2 \quad (\text{B.48})$$

É interessante comparar estes resultados com os conhecidos do caso isotrópico, onde:

$$\varepsilon_z = \varepsilon_t = \varepsilon_r$$

$$E_r = \frac{1}{K^2} \left\{ \frac{\partial^2 E_z}{\partial r \partial z} - \frac{j\omega \mu_0}{r} \frac{\partial H_z}{\partial \phi} \right\}$$

$$E_\phi = \frac{1}{K^2} \left\{ \frac{1}{r} \frac{\partial^2 E_z}{\partial \phi \partial z} + j\omega \mu_0 \frac{\partial H_z}{\partial r} \right\}$$

$$H_r = \frac{1}{K^2} \left\{ \frac{\partial^2 H_z}{\partial r \partial z} + \frac{j\omega \varepsilon_0 \varepsilon_r}{r} \frac{\partial E_z}{\partial \phi} \right\}$$

$$H_\phi = \frac{1}{K^2} \left\{ \frac{1}{r} \frac{\partial^2 H_z}{\partial \phi \partial z} - j\omega \varepsilon_0 \varepsilon_r \frac{\partial E_z}{\partial r} \right\}$$

e

$$K^2 = k_0^2 \varepsilon_r + \gamma^2 = k_0^2 \varepsilon_r - \beta^2, \text{ para meios sem perdas } (\alpha = 0).$$

Pode-se verificar que, quando o meio é isotrópico, as equações ficam as mesmas conhecidas na literatura.

APENDICE C

Obtenção da Admitância de Superfície em $r = r_0$ no Lado da Corrugação.

Os modos transverso-magnéticos a “z” são obtidos a partir dos potenciais vetores [9]. Fazendo-se:

$$\begin{cases} \bar{\Pi}_e = \Pi_{ez} \hat{z} \\ \bar{\Pi}_h = 0 \end{cases} \quad (\text{C.1})$$

O potencial vetor $\bar{\Pi}_e$ deve satisfazer a equação de onda:

$$\nabla^2 \Pi_{ez}(r, \phi, z) + K_1^2 \Pi_{ez}(r, \phi, z) = 0 \quad (\text{C.2})$$

Na região de corrugação, a constante de propagação é zero, pois na corrugação não existe propagação. Então: $r_0 \leq r \leq r_2 \Rightarrow e^{-\gamma_z z} = e^{-0 \cdot z} = 1$.

A solução da equação (C.2) é dada por:

$$\Pi_{ez} = [C_n J_n(K_1 r) + D_n Y_n(K_1 r)] \cos(n\phi) \quad (\text{C.3})$$

Sabe-se que:

$$\varepsilon_z \rightarrow 1,$$

$$E_z = K_1^2 \Pi_e,$$

$$K_1^2 = k_0^2 + \gamma_z^2$$

então:

$$E_z(r, \phi, z) = K_1^2 [C_n J_n(K_1 r) + D_n Y_n(K_1 r)] \cos(n\phi) \quad (\text{C.4})$$

como em $r = r_2$,

$$E_z(r, \phi, z) = 0 \quad (\text{C.5})$$

logo:

$$C_n J_n(K_1 r_2) + D_n Y_n(K_1 r_2) = 0 \quad (\text{C.6})$$

resultando em:

$$D_n = -\frac{J_n(K_1 r_2)}{Y_n(K_1 r_2)} C_n \quad (\text{C.7})$$

Substituindo (C.7) na equação (C.4) tem-se:

$$E_z = K_1^2 C_n \left[J_n(K_1 r) - \frac{J_n(K_1 r_2)}{Y_n(K_1 r_2)} Y_n(K_1 r) \right] \cos(n\phi) \quad (\text{C.8})$$

considerando que não existe propagação na corrugação, ou seja, $\gamma_z = 0 \Rightarrow K_1 = k_0$, obtém-se a equação:

$$E_z = \frac{k_0^2 C_n}{Y_n(k_0 r_2)} \left[J_n(k_0 r) Y_n(k_0 r_2) - J_n(k_0 r_2) Y_n(k_0 r) \right] \cos(n\phi) \quad (\text{C.9})$$

O campo magnético na direção $\hat{\phi}$ é dado por:

$$H_\phi = \frac{1}{r} \frac{\partial^2 \Pi_h}{\partial \phi \partial z} - j\omega \epsilon_0 \epsilon_t \frac{\partial \Pi_e}{\partial r} \quad (\text{C.10})$$

Mas, $\Pi_h = 0$, então:

$$H_\phi = -j\omega \epsilon_0 \epsilon_t \frac{\partial \Pi_e}{\partial r} \quad (\text{C.11})$$

Na região de corrugação, a permissividade elétrica é igual a 1 e só existe componente na direção z, ou seja:

$$H_\phi(r, \phi, z) = \frac{-j\omega \epsilon_0}{K_1^2} \left\{ K_1^2 \frac{C_n}{Y_n(K_1 r_2)} \left[J_n'(K_1 r) Y_n(K_1 r_2) - J_n(K_1 r_2) Y_n'(K_1 r) \right] \right\} \cos(n\phi) \quad (\text{C.12})$$

ou

$$H_\phi(r, \phi, z) = -j\omega \epsilon_0 K_1 \frac{C_n}{Y_n(K_1 r_2)} \left[J_n'(K_1 r) Y_n(K_1 r_2) - J_n(K_1 r_2) Y_n'(K_1 r) \right] \cos(n\phi) \quad (\text{C.13})$$

com $K_1 = k_0$ na corrugação:

$$H_\phi(r, \phi, z) = -j\omega \epsilon_0 k_0 \frac{C_n}{Y_n(k_0 r_2)} \left[J_n'(k_0 r) Y_n(k_0 r_2) - J_n(k_0 r_2) Y_n'(k_0 r) \right] \cos(n\phi) \quad (\text{C.14})$$

A admitância de superfície é dada por:

$$Y_s(r) = \frac{H_\phi}{E_z} = \frac{-j\omega \epsilon_0 k_0}{k_0^2} \left[\frac{\left\{ \frac{C_n}{Y_n(k_0 r_2)} \left[J_n'(k_0 r) Y_n(k_0 r_2) - J_n(k_0 r_2) Y_n'(k_0 r) \right] \right\} \cos(n\phi)}{\frac{C_n}{Y_n(k_0 r_2)} \left[J_n(k_0 r) Y_n(k_0 r_2) - J_n(k_0 r_2) Y_n(k_0 r) \right] \cos(n\phi)} \right] \quad (\text{C.15})$$

ou

$$Y_S(r) = \frac{-j\omega\varepsilon_0}{k_0} \left[\frac{J'_n(k_0 r)Y_n(k_0 r_2) - J_n(k_0 r_2)Y'_n(k_0 r)}{J_n(k_0 r)Y_n(k_0 r_2) - J_n(k_0 r_2)Y_n(k_0 r)} \right] \quad (\text{C.16})$$

sabendo que $y_0 = \frac{\omega\varepsilon_0}{k_0}$, dividindo e multiplicando $Y_S(r)$ por rk_0 , tem-se:

$$Y_S(r) = -j \frac{y_0}{rk_0} S_n(k_0, r, r_2) \quad (\text{C.17})$$

com

$$S_n(k, x, y) = (kx) \left\{ \frac{J'_n(kx)Y_n(ky) - J_n(ky)Y'_n(kx)}{J_n(kx)Y_n(ky) - J_n(ky)Y_n(kx)} \right\} \quad (\text{C.18})$$

Portanto, a admitância de superfície em $r = r_0$ é dada por:

$$Y_S(r_0) = \frac{-jy_0}{k_0 r_0} S_n(k_0, r_0, r_2) \quad (\text{C.19})$$

onde

$$y_0 = \sqrt{\frac{\varepsilon_0}{\mu_0}} \text{ - admitância do espaço livre}$$

e

$$S_n(k_0, r_0, r_2) = (k_0 r_0) \left\{ \frac{J'_n(k_0 r_0)Y_n(k_0 r_2) - J_n(k_0 r_2)Y'_n(k_0 r_0)}{J_n(k_0 r_0)Y_n(k_0 r_2) - J_n(k_0 r_2)Y_n(k_0 r_0)} \right\} \quad (\text{C.20})$$

APENDICE D

Obtenção dos Coeficientes $A_n, B_n, C_n, D_n, E_n, F_n$

O sistema de equações dado pelas Equações 2.27, 2.29, 2.32, 2.34, 2.36 e 2.38 do Capítulo 2, pode ser escrito na forma matricial: $[M].[Cref] = [0]$, onde M é uma matriz 6 x 6 e Cref é uma matriz coluna. Tem-se:

$$\begin{pmatrix} \alpha_{11} & \alpha_{12} & \alpha_{13} & \alpha_{14} & \alpha_{15} & \alpha_{16} \\ \alpha_{21} & \alpha_{22} & \alpha_{23} & \alpha_{24} & \alpha_{25} & \alpha_{26} \\ \alpha_{31} & \alpha_{32} & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} \\ \alpha_{41} & \alpha_{42} & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} \\ \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} \\ \alpha_{61} & \alpha_{62} & \alpha_{63} & \alpha_{64} & \alpha_{65} & \alpha_{66} \end{pmatrix} \begin{pmatrix} A_n \\ B_n \\ C_n \\ D_n \\ E_n \\ F_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad (D.1)$$

onde:

$$\alpha_{22} = \alpha_{25} = \alpha_{26} = \alpha_{31} = \alpha_{32} = \alpha_{41} = \alpha_{42} = \alpha_{61} = \alpha_{63} = \alpha_{64} = 0 \quad (D.2)$$

$$\alpha_{11} = \frac{n\gamma}{K^2 r_1} J_n(Kr_1) \quad (D.3)$$

$$\alpha_{12} = \frac{j\omega\mu_0}{K} J_n'(Kr_1) \quad (D.4)$$

$$\alpha_{13} = -\frac{n\gamma}{K_1^2 r_1} J_n(K_1 r_1) \quad (D.5)$$

$$\alpha_{14} = -\frac{n\gamma}{K_1^2 r_1} Y_n(K_1 r_1) \quad (D.6)$$

$$\alpha_{15} = -\frac{j\omega\mu_0}{K_1} J_n'(K_1 r_1) \quad (D.7)$$

$$\alpha_{16} = -\frac{j\omega\mu_0}{K_1} Y_n'(K_1 r_1) \quad (D.8)$$

$$\alpha_{21} = J_n(Kr_1) \quad (D.9)$$

$$\alpha_{23} = -J_n(K_1 r_1) \quad (D.10)$$

$$\alpha_{24} = -Y_n(K_1 r_1) \quad (D.11)$$

$$\alpha_{33} = K_1^2 J_n(K_1 r_0) Y_s(r_0) + j\omega\varepsilon_0 K_1 J_n'(K_1 r_0) \quad (D.12)$$

$$\alpha_{34} = K_1^2 Y_n(K_1 r_0) Y_s(r_0) + j\omega\varepsilon_0 K_1 Y_n'(K_1 r_0) \quad (D.13)$$

$$\alpha_{35} = \frac{n\gamma}{r_0} J_n(K_1 r_0) \quad (D.14)$$

$$\alpha_{36} = \frac{n\gamma}{r_0} Y_n(K_1 r_0) \quad (D.15)$$

$$\alpha_{43} = \frac{n\gamma}{r_0} J_n(K_1 r_0) \quad (D.16)$$

$$\alpha_{44} = \frac{n\gamma}{r_0} Y_n(K_1 r_0) \quad (D.17)$$

$$\alpha_{45} = j\omega\mu_0 K_1 J_n'(K_1 r_0) \quad (D.18)$$

$$\alpha_{46} = j\omega\mu_0 K_1 Y_n'(K_1 r_0) \quad (D.19)$$

$$\alpha_{51} = -\frac{j\omega\varepsilon_0\varepsilon_{ti}}{K} J_n'(K r_1) \quad (D.20)$$

$$\alpha_{52} = -\frac{n\gamma\varepsilon_{ti}}{K^2 r_1} J_n(K r_1) \quad (D.21)$$

$$\alpha_{53} = \frac{j\omega\varepsilon_0}{K_1} J_n'(K_1 r_1) \quad (D.22)$$

$$\alpha_{54} = \frac{j\omega\varepsilon_0}{K_1} Y_n'(K_1 r_1) \quad (D.23)$$

$$\alpha_{55} = \frac{n\gamma}{r_1 K_1^2} J_n(K_1 r_1) \quad (D.24)$$

$$\alpha_{56} = \frac{n\gamma}{r_1 K_1^2} Y_n(K_1 r_1) \quad (D.25)$$

$$\alpha_{62} = J_n(K r_1) \quad (D.26)$$

$$\alpha_{65} = -J_n(K_1 r_1) \quad (D.27)$$

$$\alpha_{66} = -Y_n(K_1 r_1) \quad (D.28)$$

onde, $\gamma = \alpha + j\beta \cong j\beta = \gamma_z =$ constante de propagação na direção z para meios sem perdas, e $Y_s(r_0)$ é a admitância de superfície em $r = r_0$, dada pela equação C.17 (apêndice C).

Sabendo que:

$$k_0^2 = \omega^2 \mu_0 \varepsilon_0 \quad (D.29)$$

$$\omega\varepsilon_0 = k_0 y_o \quad (D.30)$$

$$\omega\mu_0 = \frac{k_0}{y_o} \quad (D.31)$$

E, substituindo D.29, D.30 e D.31 nas Equações D.3 a D.28, pode-se reescrever a matriz como:

$$\begin{pmatrix} \alpha_{11} & \alpha_{12} & \alpha_{13} & \alpha_{14} & \alpha_{15} & \alpha_{16} \\ \alpha_{21} & 0 & \alpha_{23} & \alpha_{24} & 0 & 0 \\ 0 & 0 & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} \\ \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} \end{pmatrix} \begin{pmatrix} A_n \\ B_n \\ C_n \\ D_n \\ E_n \\ F_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad (\text{D.32})$$

com

$$\alpha_{22} = \alpha_{25} = \alpha_{26} = \alpha_{31} = \alpha_{32} = \alpha_{41} = \alpha_{42} = \alpha_{61} = \alpha_{63} = \alpha_{64} = 0 \quad (\text{D.33})$$

$$\alpha_{11} = \frac{j\beta n}{K^2 r_1} J_n(Kr_1) \quad (\text{D.34})$$

$$\alpha_{12} = \frac{jk_0}{Ky_0} J'_n(Kr_1) \quad (\text{D.35})$$

$$\alpha_{13} = -\frac{j\beta n}{K_1^2 r_1} J_n(K_1 r_1) \quad (\text{D.36})$$

$$\alpha_{14} = -\frac{j\beta n}{K_1^2 r_1} Y_n(K_1 r_1) \quad (\text{D.37})$$

$$\alpha_{15} = -\frac{jk_0}{K_1 y_0} J'_n(K_1 r_1) \quad (\text{D.38})$$

$$\alpha_{16} = -\frac{jk_0}{K_1 y_0} Y'_n(K_1 r_1) \quad (\text{D.39})$$

$$\alpha_{21} = J_n(Kr_1) \quad (\text{D.40})$$

$$\alpha_{23} = -J_n(K_1 r_1) \quad (\text{D.41})$$

$$\alpha_{24} = -Y_n(K_1 r_1) \quad (\text{D.42})$$

$$\alpha_{33} = K_1^2 J_n(K_1 r_0) Y_s(r_0) + jk_0 y_0 K_1 J'_n(K_1 r_0) \quad (\text{D.43})$$

$$\alpha_{34} = K_1^2 Y_n(K_1 r_0) Y_s(r_0) + jk_0 y_0 K_1 Y'_n(K_1 r_0) \quad (\text{D.44})$$

$$\alpha_{35} = \frac{j\beta n}{r_0} J_n(K_1 r_0) \quad (\text{D.45})$$

$$\alpha_{36} = \frac{j\beta n}{r_0} Y_n(K_1 r_0) \quad (\text{D.46})$$

$$\alpha_{43} = \frac{j\beta n}{r_0} J_n(K_1 r_0) \quad (\text{D.47})$$

$$\alpha_{44} = \frac{j\beta n}{r_0} Y_n(K_1 r_0) \quad (D.48)$$

$$\alpha_{45} = j \frac{k_0}{y_0} K_1 J'_n(K_1 r_0) \quad (D.49)$$

$$\alpha_{46} = j \frac{k_0}{y_0} K_1 Y'_n(K_1 r_0) \quad (D.50)$$

$$\alpha_{51} = -\frac{jk_0 y_0 \varepsilon_{ii}}{K} J'_n(K r_1) \quad (D.51)$$

$$\alpha_{52} = -\frac{j\beta n}{K^2 r_1^2} \varepsilon_{ii} J_n(K r_1) \quad (D.52)$$

$$\alpha_{53} = \frac{jk_0 y_0}{K_1} J'_n(K_1 r_1) \quad (D.53)$$

$$\alpha_{54} = \frac{jk_0 y_0}{K_1} Y'_n(K_1 r_1) \quad (D.54)$$

$$\alpha_{55} = \frac{j\beta n}{r_1 K_1^2} J_n(K_1 r_1) \quad (D.55)$$

$$\alpha_{56} = \frac{j\beta n}{r_1 K_1^2} Y_n(K_1 r_1) \quad (D.56)$$

$$\alpha_{62} = J_n(K r_1) \quad (D.57)$$

$$\alpha_{65} = -J_n(K_1 r_1) \quad (D.58)$$

$$\alpha_{66} = -Y_n(K_1 r_1) \quad (D.59)$$

Gerando a matriz aumentada e triangularizando a matriz, ou seja, zerando os valores abaixo da diagonal principal [13] obtém-se o determinante da matriz. Para isso, usando a matriz aumentada:

$$\begin{pmatrix} \alpha_{11} & \alpha_{12} & \alpha_{13} & \alpha_{14} & \alpha_{15} & \alpha_{16} & 0 \\ \alpha_{21} & 0 & \alpha_{23} & \alpha_{24} & 0 & 0 & 0 \\ 0 & 0 & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \end{pmatrix} \quad (D.60)$$

Trocando algumas linhas para facilitar a triangularização, escreve-se a nova matriz:

$$\begin{pmatrix} \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ \alpha_{21} & 0 & \alpha_{23} & \alpha_{24} & 0 & 0 & 0 \\ \alpha_{11} & \alpha_{12} & \alpha_{13} & \alpha_{14} & \alpha_{15} & \alpha_{16} & 0 \\ 0 & 0 & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} & 0 \end{pmatrix} \quad (\text{D.61})$$

Zerando a coluna 1 das linhas 4 e 5 com a linha 1 (α_{5n}):

$$\begin{pmatrix} \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ 0 & -\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} & \alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} & \alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} & -\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} & -\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} & 0 \\ 0 & \alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} & \alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} & \alpha_{14} - \alpha_{11} \frac{\alpha_{54}}{\alpha_{51}} & \alpha_{15} - \alpha_{11} \frac{\alpha_{55}}{\alpha_{51}} & \alpha_{16} - \alpha_{11} \frac{\alpha_{56}}{\alpha_{51}} & 0 \\ 0 & 0 & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} & 0 \end{pmatrix} \quad (\text{D.62})$$

Zerando a coluna 2 das linhas 4 e 5 a partir da linha 2 (α_{6n}):

$$\begin{pmatrix} \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ 0 & 0 & \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) & \left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) & \begin{pmatrix} -\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} + \\ \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \alpha_{65} \end{pmatrix} & \begin{pmatrix} -\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} + \\ \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \alpha_{66} \end{pmatrix} & 0 \\ 0 & 0 & \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) & \left(\alpha_{14} - \alpha_{11} \frac{\alpha_{54}}{\alpha_{51}} \right) & \begin{pmatrix} \alpha_{15} - \alpha_{11} \frac{\alpha_{55}}{\alpha_{51}} - \\ \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \alpha_{65} \end{pmatrix} & \begin{pmatrix} \alpha_{16} - \alpha_{11} \frac{\alpha_{56}}{\alpha_{51}} - \\ \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \alpha_{66} \end{pmatrix} & 0 \\ 0 & 0 & \alpha_{33} & \alpha_{34} & \alpha_{35} & \alpha_{36} & 0 \end{pmatrix} \quad (\text{D.63})$$

Zerando a coluna 3 das linhas 4, 5 e 6 a partir da linha 3 (α_{4n}):

$$\begin{pmatrix} \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ 0 & 0 & 0 & L44n & L45n & L46n & 0 \\ 0 & 0 & 0 & L54n & L55n & L56n & 0 \\ 0 & 0 & 0 & \alpha_{34} \frac{\alpha_{44}}{\alpha_{43}} & \alpha_{35} \frac{\alpha_{45}}{\alpha_{43}} & \alpha_{36} \frac{\alpha_{46}}{\alpha_{43}} & 0 \end{pmatrix} \quad (\text{D.64})$$

com

$$L44n = \left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \quad (D.65)$$

$$L45n = -\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} \quad (D.66)$$

$$L46n = -\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} \quad (D.67)$$

$$L54n = \alpha_{14} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \quad (D.68)$$

$$L55n = \alpha_{15} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} \quad (D.69)$$

$$L56n = \alpha_{16} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} \quad (D.70)$$

Zerando a coluna 4 das linhas 5 e 6 com linha 4:

$$\begin{pmatrix} \alpha_{51} & \alpha_{52} & \alpha_{53} & \alpha_{54} & \alpha_{55} & \alpha_{56} & 0 \\ 0 & \alpha_{62} & 0 & 0 & \alpha_{65} & \alpha_{66} & 0 \\ 0 & 0 & \alpha_{43} & \alpha_{44} & \alpha_{45} & \alpha_{46} & 0 \\ 0 & 0 & 0 & L44n & L45n & L46n & 0 \\ 0 & 0 & 0 & 0 & L55n & L56n & 0 \\ 0 & 0 & 0 & 0 & L65n & L66n & 0 \end{pmatrix} \quad (D.71)$$

$$L44n = \alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} + \alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \frac{\alpha_{64}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \quad (D.72)$$

$$L45n = -\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} \quad (D.73)$$

$$L46n = -\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} \quad (D.74)$$

$$L55n = \alpha_{15} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} - \frac{\left(\alpha_{14} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \right) \left(-\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} \right)}{\left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}}} \quad (D.75)$$

$$L565n = \alpha_{16} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{12} - \alpha_{11} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} -$$

$$\left(\alpha_{14} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} - \left(\alpha_{13} - \alpha_{11} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \right) \frac{\left(-\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} \right)}{\left(\left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \right)} \quad (D.76)$$

$$L65n = \alpha_{35} \frac{\alpha_{45}}{\alpha_{43}} - \left(\alpha_{34} \frac{\alpha_{44}}{\alpha_{43}} \right) \frac{\left(-\alpha_{21} \frac{\alpha_{55}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{65}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{45}}{\alpha_{43}} \right)}{\left(\left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \right)} \quad (D.77)$$

$$L66n = \alpha_{36} \frac{\alpha_{46}}{\alpha_{43}} \left(\alpha_{34} \frac{\alpha_{44}}{\alpha_{43}} \right) \frac{\left(-\alpha_{21} \frac{\alpha_{56}}{\alpha_{51}} + \left(\alpha_{21} \frac{\alpha_{52}}{\alpha_{51}} \right) \frac{\alpha_{66}}{\alpha_{62}} - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{46}}{\alpha_{43}} \right)}{\left(\left(\alpha_{24} - \alpha_{21} \frac{\alpha_{54}}{\alpha_{51}} \right) - \left(\alpha_{23} - \alpha_{21} \frac{\alpha_{53}}{\alpha_{51}} \right) \frac{\alpha_{44}}{\alpha_{43}} \right)} \quad (D.78)$$

Usando a matriz resultante da Equação (D.71) para calcular o determinante, tem-se os coeficientes em função de An (normalizando em An=1).

Da linha 5:

$$L555n.An + L556n.Cn = 0 \Rightarrow$$

$$Cn = -\frac{L555n}{L556n} An \quad (D.79)$$

ou

$$Cn = NAn \quad (D.80)$$

com

$$N = -\frac{L555n}{L556n} \quad (D.81)$$

Da linha 4:

$$L44nBn + L45nAn + L46nCn = 0 \Rightarrow$$

$$L44nBn + L45nAn - L46n \frac{L555}{L556} An = 0 \quad (D.82)$$

$$Bn = -\frac{\left(L45n + \frac{L46nL555}{L556} \right)}{L44n} An$$

ou

$$Bn = MAn \quad (D.83)$$

com

$$M = -\frac{\left(L45n + \frac{L46nL555}{L556}\right)}{L44n} \quad (D.84)$$

Da linha 3:

$$\alpha_{43}Dn + \alpha_{44}Bn + \alpha_{45}An + \alpha_{46}Cn = 0 \Rightarrow$$

$$\alpha_{43}Dn - \alpha_{44}\left(\frac{\left(L45n + \frac{L46nL555}{L556}\right)}{L44n}\right)An + \alpha_{45}An - \alpha_{46}\left(\frac{L555n}{L556n}An\right) = 0 \quad (D.85)$$

$$Dn = \frac{\alpha_{44}\left(\frac{\left(L45n + \frac{L46nL555}{L556}\right)}{L44n}\right) - \alpha_{45} + \alpha_{46}\left(\frac{L555n}{L556n}\right)}{\alpha_{43}}An$$

ou

$$Dn = PAn \quad (D.86)$$

com

$$P = \frac{\alpha_{44}\left(\frac{\left(L45n + \frac{L46nL555}{L556}\right)}{L44n}\right) - \alpha_{45} + \alpha_{46}\left(\frac{L555n}{L556n}\right)}{\alpha_{43}} \quad (D.87)$$

Da linha 2:

$$\alpha_{62}Fn + \alpha_{65}An + \alpha_{66}Cn = 0 \Rightarrow \alpha_{62}Fn + \alpha_{65}An - \alpha_{66}\left(\frac{L555n}{L556n}An\right) = 0$$

$$\Rightarrow Fn = \frac{\alpha_{66}\left(\frac{L555n}{L556n}\right) - \alpha_{65}}{\alpha_{62}}An \quad (D.88)$$

ou

$$Fn = LAn \quad (D.89)$$

com

$$L = \frac{\alpha_{66} \left(\frac{L555n}{L556n} \right) - \alpha_{65}}{\alpha_{62}} \quad (D.90)$$

Da linha 1:

$$\alpha_{51}En + \alpha_{52}Fn + \alpha_{53}Dn + \alpha_{54}Bn + \alpha_{55}An + \alpha_{56}Cn = 0 \Rightarrow$$

$$\alpha_{51}En + \alpha_{52} \left(\frac{\alpha_{66} \left(\frac{L555n}{L556n} \right) - \alpha_{65}}{\alpha_{62}} \right) An + \alpha_{53} \left(\frac{\alpha_{44} \left(\frac{\left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) - \alpha_{45} + \alpha_{46} \left(\frac{L555n}{L556n} \right)}{\alpha_{43}} \right) An - \left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) An + \alpha_{55}An - \alpha_{56} \frac{L555n}{L556n} An = 0 \Rightarrow \quad (D.91)$$

logo:

$$En = \frac{An}{\alpha_{51}} \left\{ -\alpha_{52} \left(\frac{\alpha_{66} \left(\frac{L555n}{L556n} \right) - \alpha_{65}}{\alpha_{62}} \right) - \alpha_{53} \left(\frac{\alpha_{44} \left(\frac{\left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) - \alpha_{45} + \alpha_{46} \left(\frac{L555n}{L556n} \right)}{\alpha_{43}} \right) + \left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) - \alpha_{55} + \alpha_{56} \frac{L555n}{L556n} \right\} \quad (D.92)$$

ou

$$En = QAn \quad (D.93)$$

com

$$Q = \left\{ \begin{array}{l} -\alpha_{52} \left(\frac{\alpha_{66} \left(\frac{L555n}{L556n} \right) - \alpha_{65}}{\alpha_{62}} \right) - \alpha_{53} \left(\frac{\alpha_{44} \left(\frac{\left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) - \alpha_{45} + \alpha_{46} \left(\frac{L555n}{L556n} \right)}{\alpha_{43}} \right) + \left(\frac{L45n + \frac{L46nL555}{L556}}{L44n} \right) - \alpha_{55} + \alpha_{56} \frac{L555n}{L556n} \end{array} \right\} \frac{1}{\alpha_{51}} \quad (D.94)$$

APÊNDICE E

Desenvolvimento de Integrais de Campos usadas [10,14]

E.1

Cos(nφ)sen(φ-φ')

$$\begin{aligned}
 & \int_0^{2\pi} \cos(n\phi') \operatorname{sen}(\phi - \phi') e^{jkr' \operatorname{sen} \theta \cos(\phi - \phi')} d\phi' = \\
 & \alpha = \phi - \phi' \Rightarrow \phi' = \phi - \alpha \\
 & d\alpha = -d\phi' \\
 & = -\int_{\phi}^{\phi-2\pi} \cos n(\phi - \alpha) \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
 & -\int_{\phi}^{\phi-2\pi} \cos n\phi \cos n\alpha \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \\
 & \int_{\phi}^{\phi-2\pi} \sin n\phi \sin(n\alpha) \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
 & = -\cos n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} + e^{-jn\alpha})}{2} \frac{(e^{j\alpha} - e^{-j\alpha})}{2j} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \\
 & \operatorname{sen} n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} - e^{-jn\alpha})}{2j} \frac{(e^{j\alpha} - e^{-j\alpha})}{2j} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
 & = -\frac{\cos n\phi}{4j} \int_0^{2\pi} \frac{(e^{j\alpha(n+1)} + e^{-j\alpha(n-1)})}{-e^{j\alpha(n-1)} - e^{-j\alpha(n+1)}} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \\
 & \frac{\operatorname{sen} n\phi}{4} \int_0^{2\pi} \frac{(e^{j\alpha(n+1)} - e^{-j\alpha(n-1)})}{-e^{j\alpha(n-1)} + e^{-j\alpha(n+1)}} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
 & = -\frac{\cos n\phi}{4j} \left\{ \begin{aligned} & \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) + \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \operatorname{sen} \theta) \\ & - \frac{2\pi}{j^{-(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) - \frac{2\pi}{j^{(n+1)}} J_{-(n+1)}(kr' \operatorname{sen} \theta) \end{aligned} \right\} - \\
 & \frac{\operatorname{sen} n\phi}{4} \left\{ \begin{aligned} & \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) - \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \operatorname{sen} \theta) \\ & - \frac{2\pi}{j^{-(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) + \frac{2\pi}{j^{(n+1)}} J_{-(n+1)}(kr' \operatorname{sen} \theta) \end{aligned} \right\} = \\
 & = -\frac{2\pi \cos n\phi}{4j} \left\{ \begin{aligned} & \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) + \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) \\ & - \frac{1}{j^{-(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) - \frac{(-1)^{(n+1)}}{j^{(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) \end{aligned} \right\} - \\
 & \frac{2\pi \operatorname{sen} n\phi}{4} \left\{ \begin{aligned} & \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) - \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) \\ & - \frac{1}{j^{-(n-1)}} J_{(n-1)}(kr' \operatorname{sen} \theta) + \frac{(-1)^{(n+1)}}{j^{(n+1)}} J_{(n+1)}(kr' \operatorname{sen} \theta) \end{aligned} \right\} = \\
 & = -\frac{\pi \cos n\phi}{2j} \left\{ \left(\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \operatorname{sen} \theta) + \left(\frac{(-1)^{(n-1)}}{j^{(n-1)}} - \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \operatorname{sen} \theta) \right\} \\
 & - \frac{\pi \operatorname{sen} n\phi}{2} \left\{ \left(\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \operatorname{sen} \theta) - \left(\frac{(-1)^{(n-1)}}{j^{(n-1)}} + \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \operatorname{sen} \theta) \right\}
 \end{aligned}$$

OU

$$\begin{aligned}
& \int_0^{2\pi} \cos(n\phi') \sin(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = -\frac{\pi \cos n\phi}{2j} \left\{ \left(\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(\frac{(-1)^{(n-1)}}{j^{(n-1)}} - \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} + \\
& + \frac{\pi \sin n\phi}{2} \left\{ \left(\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) - \left(\frac{(-1)^{(n-1)}}{j^{(n-1)}} + \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} = \\
& = \frac{\pi}{2} \left\{ -\cos n\phi \left[\left(\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(-\frac{(-1)^{(n-1)}}{j^{(n-1)}} - \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right] + \right. \\
& \left. \sin n\phi \left[\left(\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) - \left(-\frac{(-1)^{(n-1)}}{j^{(n-1)}} + \frac{1}{j^{-(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right] \right\} = \\
& \frac{\pi}{2} \left\{ -\cos n\phi (j^{(n)} + j^{(n-2)}) \{ J_{(n+1)}(kr' \text{sen } \theta) - J_{(n-1)}(kr' \text{sen } \theta) \} + \sin n\phi (j^{(n+1)} - j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen } \theta) + J_{(n-1)}(kr' \text{sen } \theta) \} \right\} = \\
& = \frac{\pi}{2} \left\{ -\cos n\phi (j^{(n)} - j^{(n)}) \{ J_{(n+1)}(kr' \text{sen } \theta) - J_{(n-1)}(kr' \text{sen } \theta) \} + \sin n\phi (j^{(n+1)} - j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen } \theta) + J_{(n-1)}(kr' \text{sen } \theta) \} \right\} = \\
& \frac{\pi}{2} \sin n\phi (j^{(n+1)} - j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen } \theta) + J_{(n-1)}(kr' \text{sen } \theta) \} = \frac{\pi}{2} \sin n\phi j^{(n)} \left(\frac{-1-1}{j} \right) \{ J_{(n+1)}(kr' \text{sen } \theta) + J_{(n-1)}(kr' \text{sen } \theta) \} = \\
& -\pi \sin n\phi j^{(n-1)} \{ J_{(n+1)}(kr' \text{sen } \theta) + J_{(n-1)}(kr' \text{sen } \theta) \} = -\frac{2\pi \sin n\phi j^{(n-1)}}{k \sin \theta} J_{(n)}(kr' \text{sen } \theta)
\end{aligned}$$

Então

$$\int_0^{2\pi} \cos(n\phi') \sin(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = -\frac{2\pi \sin n\phi j^{(n-1)}}{k \sin \theta} J_{(n)}(kr' \text{sen } \theta) \quad (\text{E.1})$$

E.2

Cos(nφ)cos(φ-φ')

$$\begin{aligned}
& \int_0^{2\pi} \cos(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \\
& \alpha = \phi - \phi' \Rightarrow \phi' = \phi - \alpha \\
& d\alpha = -d\phi' \\
& = -\int_{\phi}^{\phi-2\pi} \cos n(\phi - \alpha) \cos \alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = -\int_{\phi}^{\phi-2\pi} \cos n\phi \cos n\alpha \cos \alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha - \int_{\phi}^{\phi-2\pi} \sin n\phi \sin(n\alpha) \cos \alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = \\
& = \cos n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} + e^{-jn\alpha})}{2} \frac{(e^{j\alpha} + e^{-j\alpha})}{2} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha - \sin n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} - e^{-jn\alpha})}{2j} \frac{(e^{j\alpha} + e^{-j\alpha})}{2} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = \\
& = \frac{\cos n\phi}{4} \int_0^{2\pi} (e^{jn\alpha} + e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha}) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha - \frac{\sin n\phi}{4j} \int_0^{2\pi} (e^{jn\alpha} - e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha}) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = \\
& = \frac{\cos n\phi}{4} \int_0^{2\pi} (e^{j\alpha(n+1)} + e^{-j\alpha(n-1)} + e^{j\alpha(n-1)} + e^{-j\alpha(n+1)}) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha - \frac{\sin n\phi}{4j} \int_0^{2\pi} (e^{j\alpha(n+1)} - e^{-j\alpha(n-1)} + e^{j\alpha(n-1)} - e^{-j\alpha(n+1)}) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = \\
& = \frac{\cos n\phi}{4} \left\{ \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) + \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen } \theta) + \frac{2\pi}{j^{-(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) + \frac{2\pi}{j^{(n+1)}} J_{-(n+1)}(kr' \text{sen } \theta) \right\} \\
& - \frac{\sin n\phi}{4j} \left\{ \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) - \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen } \theta) + \frac{2\pi}{j^{-(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) - \frac{2\pi}{j^{(n+1)}} J_{-(n+1)}(kr' \text{sen } \theta) \right\} = \\
& = \frac{2\pi \cos n\phi}{4} \left\{ \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) + \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) + \frac{1}{j^{-(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) + \frac{(-1)^{(n+1)}}{j^{(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) \right\} \\
& - \frac{2\pi \sin n\phi}{4j} \left\{ \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) - \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) + \frac{1}{j^{-(n-1)}} J_{(n-1)}(kr' \text{sen } \theta) - \frac{(-1)^{(n+1)}}{j^{(n+1)}} J_{(n+1)}(kr' \text{sen } \theta) \right\} =
\end{aligned}$$

OU

$$\begin{aligned}
&= \frac{\pi \cos n\phi}{2} \left\{ \left(\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(\frac{1}{j^{-(n-1)}} + \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} \\
&- \frac{\pi \sin n\phi}{2j} \left\{ \left(\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(\frac{1}{j^{-(n-1)}} - \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} = \\
&= \frac{\pi \cos n\phi}{2} \left\{ \left(j^{(n+1)} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(j^{(n-1)} - \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} + \\
&- \frac{\pi \sin n\phi}{2} \left\{ \left(j^n + \frac{(-1)^{(n+1)}}{j^{(n+2)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(j^{(n-2)} + \frac{(-1)^{(n-1)}}{j^n} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} = \\
&= \frac{\pi \cos n\phi}{2} \left\{ \left(j^{(n+1)} - \frac{(j)^{(2n)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(j^{(n-1)} - \frac{(j)^{(2n)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} + \\
&- \frac{\pi \sin n\phi}{2} \left\{ \left(j^n + \frac{(j)^{(2n)}}{j^{(n+2)}} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(j^{(n-2)} + \frac{(j)^{(2n)}}{j^n} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} = \\
&= \frac{\pi \cos n\phi}{2} \left\{ \left(j^{(n+1)} - j^{(n-1)} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(j^{(n-1)} - j^{(n+1)} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} + \\
&- \frac{\pi \sin n\phi}{2} \left\{ \left(j^n - j^{(n)} \right) J_{(n+1)}(kr' \text{sen } \theta) + \left(-j^{(n)} + j^{(n)} \right) J_{(n-1)}(kr' \text{sen } \theta) \right\} = \\
&= \frac{\pi \cos n\phi}{2} \left(j^{(n+1)} - j^{(n-1)} \right) \left\{ J_{(n+1)}(kr' \text{sen } \theta) - J_{(n-1)}(kr' \text{sen } \theta) \right\} = -\pi \cos n\phi \left(j^{(n+1)} - j^{(n-1)} \right) J'_{(n)}(kr' \text{sen } \theta) = \\
&= 2\pi \cos n\phi j^{(n-1)} J'_{(n)}(kr' \text{sen } \theta)
\end{aligned}$$

logo:

$$\int_0^{2\pi} \cos(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = 2\pi \cos n\phi j^{(n-1)} J'_n(kr' \text{sen } \theta) \quad (\text{E.2})$$

E.3 Sen(nφ')sen(φ-φ')

$$\int_0^{2\pi} \text{sen } n\phi' \text{sen}(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = ?$$

$$\alpha = \phi - \phi' \Rightarrow \phi' = \phi - \alpha$$

$$d\alpha = -d\phi'$$

$$x = kr' \text{sen } \theta$$

$$= - \int_{\phi}^{\phi-2\pi} \text{sen}(n(\phi - \alpha)) \text{sen } \alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha =$$

$$\text{sen}(a - b) = \text{sen } a \cos b - \cos a \text{sen } b \Rightarrow$$

$$= - \int_{\phi}^{\phi-2\pi} \text{sen}(n\phi) \cos(n\alpha) \text{sen } \alpha e^{jx \cos \alpha} d\alpha + \int_{\phi}^{\phi-2\pi} \cos n\phi \text{sen}(n\alpha) \text{sen } \alpha e^{jx \cos \alpha} d\alpha =$$

$$= - \text{sen}(n\phi) \int_{\phi}^{\phi-2\pi} \cos(n\alpha) \text{sen } \alpha e^{jx \cos \alpha} d\alpha + \cos(n\phi) \int_{\phi}^{\phi-2\pi} \text{sen}(n\alpha) \text{sen } \alpha e^{jx \cos \alpha} d\alpha =$$

mas

$$\begin{aligned}\cos(n\alpha) \operatorname{sen} \alpha &= \frac{1}{2} \{ \operatorname{sen}(n\alpha + \alpha) - \operatorname{sen}(n\alpha - \alpha) \} = \frac{1}{2} \left\{ \frac{(e^{j\alpha(n+1)} - e^{-j\alpha(n+1)})}{2j} - \frac{(e^{j\alpha(n-1)} - e^{-j\alpha(n-1)})}{2j} \right\} \\ \operatorname{sen}(n\alpha) \operatorname{sen} \alpha &= \frac{1}{2} \{ -\cos(n\alpha + \alpha) + \cos(n\alpha - \alpha) \} = \frac{1}{2} \left\{ -\frac{(e^{j\alpha(n+1)} + e^{-j\alpha(n+1)})}{2} + \frac{(e^{j\alpha(n-1)} + e^{-j\alpha(n-1)})}{2} \right\}\end{aligned}$$

então

$$\begin{aligned}\int_0^{2\pi} \operatorname{sen} n\phi' \operatorname{sen}(\phi - \phi') e^{jkr' \operatorname{sen} \theta \cos(\phi - \phi')} d\phi' &= -\frac{\pi}{2j} \operatorname{sen}(n\phi) \left(\frac{1}{j^{-(n+1)}} J_{(n+1)}(u) - \frac{1}{j^{(n+1)}} J_{-(n+1)}(u) - \right. \\ &\quad \left. + \frac{1}{j^{-(n-1)}} J_{(n-1)}(u) + \frac{1}{j^{(n-1)}} J_{-(n-1)}(u) \right) + \\ &\quad + \frac{\pi}{2} \cos(n\phi) \left(-\frac{1}{j^{-(n+1)}} J_{(n+1)}(u) - \frac{1}{j^{(n+1)}} J_{-(n+1)}(u) + \right. \\ &\quad \left. + \frac{1}{j^{-(n-1)}} J_{(n-1)}(u) + \frac{1}{j^{(n-1)}} J_{-(n-1)}(u) \right)\end{aligned}$$

mas

$$J_{-(n+1)}(u) = (-1)^{(n+1)} J_{(n+1)}(u)$$

logo

$$\begin{aligned}\int_0^{2\pi} \operatorname{sen} n\phi' \operatorname{sen}(\phi - \phi') e^{jkr' \operatorname{sen} \theta \cos(\phi - \phi')} d\phi' &= \frac{\pi}{2} \left\{ \frac{1}{j} \operatorname{sen}(n\phi) \left[\left(-\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(u) + \right. \right. \\ &\quad \left. \left. + \left(\frac{1}{j^{-(n-1)}} - \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(u) \right] + \right. \\ &\quad \left. + \cos(n\phi) \left[\left(-\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(u) + \right. \right. \\ &\quad \left. \left. + \left(\frac{1}{j^{-(n-1)}} + \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(u) \right] \right\} = \\ &= \frac{\pi}{2} \left\{ \frac{1}{j} \operatorname{sen}(n\phi) \left[\left(-\frac{1}{j^{-(n+1)}} + \frac{j^{2(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(u) + \left(\frac{1}{j^{-(n-1)}} - \frac{j^{2(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(u) \right] + \right. \\ &\quad \left. + \cos(n\phi) \left[\left(-\frac{1}{j^{-(n+1)}} - \frac{j^{2(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(u) + \left(\frac{1}{j^{-(n-1)}} + \frac{j^{2(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(u) \right] \right\} = \\ &= \frac{\pi}{2} \left\{ \frac{1}{j} \operatorname{sen}(n\phi) \left[(-j^{(n+1)} + j^{(n+1)}) J_{(n+1)}(u) + (j^{(n-1)} - j^{(n-1)}) J_{(n-1)}(u) \right] + \right. \\ &\quad \left. + \cos(n\phi) \left[(-j^{(n+1)} - j^{(n+1)}) J_{(n+1)}(u) + (j^{(n-1)} + j^{(n-1)}) J_{(n-1)}(u) \right] \right\} = \\ &= \frac{\pi}{2} \left\{ \cos(n\phi) \left[(-j^{(n+1)} - j^{(n+1)}) J_{(n+1)}(u) + (j^{(n-1)} + j^{(n-1)}) J_{(n-1)}(u) \right] \right\} = \\ &= \frac{\pi}{2} \left\{ \cos(n\phi) \left[-2j^{(n+1)} J_{(n+1)}(u) + 2j^{(n-1)} J_{(n-1)}(u) \right] \right\} = \\ &= \pi j^n \cos(n\phi) \left[-j J_{(n+1)}(u) + j^{-1} J_{(n-1)}(u) \right] = \\ &= \pi j^{n-1} \cos(n\phi) \left[J_{(n+1)}(u) + J_{(n-1)}(u) \right] = \\ &= 2\pi j^{n-1} \cos(n\phi) \frac{J_n(u)}{k \operatorname{sen} \theta}\end{aligned}$$

Portanto:

$$\int_0^{2\pi} \text{sen}(n\phi') \text{sen}(\phi - \phi') e^{jkr' \text{sen} \theta \cos(\phi - \phi')} d\phi' = 2\pi j^{n-1} \cos(n\theta) \frac{J_n(kr' \text{sen} \theta)}{k \text{sen} \theta} \quad (\text{E.3})$$

E.4

Sen(n\phi)cos(\phi-\phi')

$$\begin{aligned} & \int_0^{2\pi} \text{sen}(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen} \theta \cos(\phi - \phi')} d\phi' = \\ & \alpha = \phi - \phi' \Rightarrow \phi' = \phi - \alpha \\ & d\alpha = -d\phi' \\ & = -\int_{\phi}^{\phi-2\pi} \text{sen} n(\phi - \alpha) \cos \alpha e^{jkr' \text{sen} \theta \cos \alpha} d\alpha = -\int_{\phi}^{\phi-2\pi} \text{sen} n\phi \cos n\alpha \cos \alpha e^{jkr' \text{sen} \theta \cos \alpha} d\alpha + \int_{\phi}^{\phi-2\pi} \cos n\phi \text{sen}(n\alpha) \cos \alpha e^{jkr' \text{sen} \theta \cos \alpha} d\alpha = \\ & = -\text{sen} n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} + e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha})}{2} e^{jkr' \text{sen} \theta \cos \alpha} d\alpha + \cos n\phi \int_0^{2\pi} \frac{(e^{jn\alpha} - e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha})}{2j} e^{jkr' \text{sen} \theta \cos \alpha} d\alpha = \\ & = -\frac{\text{sen} n\phi}{4} \int_0^{2\pi} (e^{jn\alpha} + e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha}) e^{jkr' \text{sen} \theta \cos \alpha} d\alpha + \frac{\cos n\phi}{4j} \int_0^{2\pi} (e^{jn\alpha} - e^{-jn\alpha})(e^{j\alpha} + e^{-j\alpha}) e^{jkr' \text{sen} \theta \cos \alpha} d\alpha = \\ & = -\frac{\text{sen} n\phi}{4} \int_0^{2\pi} (e^{j\alpha(n+1)} + e^{-j\alpha(n-1)} + e^{j\alpha(n-1)} + e^{-j\alpha(n+1)}) e^{jkr' \text{sen} \theta \cos \alpha} d\alpha + \frac{\cos n\phi}{4j} \int_0^{2\pi} (e^{j\alpha(n+1)} - e^{-j\alpha(n-1)} + e^{j\alpha(n-1)} - e^{-j\alpha(n+1)}) e^{jkr' \text{sen} \theta \cos \alpha} d\alpha = \\ & = -\frac{\text{sen} n\phi}{4} \left\{ \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen} \theta) + \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen} \theta) + \frac{2\pi}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen} \theta) + \frac{2\pi}{j^{-(n+1)}} J_{-(n+1)}(kr' \text{sen} \theta) \right\} + \\ & + \frac{\cos n\phi}{4j} \left\{ \frac{2\pi}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen} \theta) - \frac{2\pi}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen} \theta) + \frac{2\pi}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen} \theta) - \frac{2\pi}{j^{-(n+1)}} J_{-(n+1)}(kr' \text{sen} \theta) \right\} = \\ & = -\frac{2\pi \text{sen} n\phi}{4} \left\{ \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen} \theta) + \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen} \theta) + \frac{1}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen} \theta) + \frac{(-1)^{(n+1)}}{j^{-(n+1)}} J_{-(n+1)}(kr' \text{sen} \theta) \right\} + \\ & + \frac{2\pi \cos n\phi}{4} \left\{ \frac{1}{j^{-(n+1)}} J_{(n+1)}(kr' \text{sen} \theta) - \frac{(-1)^{(n-1)}}{j^{(n-1)}} J_{-(n-1)}(kr' \text{sen} \theta) + \frac{1}{j^{(n-1)}} J_{(n-1)}(kr' \text{sen} \theta) - \frac{(-1)^{(n+1)}}{j^{-(n+1)}} J_{-(n+1)}(kr' \text{sen} \theta) \right\} = \\ & = -\frac{\pi \text{sen} n\phi}{2} \left\{ \left(\frac{1}{j^{-(n+1)}} + \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{1}{j^{-(n-1)}} + \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} + \\ & + \frac{\pi \cos n\phi}{2j} \left\{ \left(\frac{1}{j^{-(n+1)}} - \frac{(-1)^{(n+1)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{1}{j^{-(n-1)}} - \frac{(-1)^{(n-1)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} = \end{aligned}$$

OU

$$\begin{aligned} & \int_0^{2\pi} \text{sen}(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen} \theta \cos(\phi - \phi')} d\phi' = -\frac{\pi \text{sen} n\phi}{2} \left\{ \left(\frac{j^{(n+1)} - (-1)^{(n)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{j^{(n-1)} - (-1)^{(n)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} + \frac{\pi \cos n\phi}{2j} \left\{ \left(\frac{j^{(n+1)} + (-1)^{(n)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{j^{(n-1)} + (-1)^{(n)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} = \\ & = -\frac{\pi \text{sen} n\phi}{2} \left\{ \left(\frac{j^{(n+1)} - (j)^{(2n)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{j^{(n-1)} - (j)^{(2n)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} + \frac{\pi \cos n\phi}{2j} \left\{ \left(\frac{j^{(n+1)} + (j)^{(2n)}}{j^{(n+1)}} \right) J_{(n+1)}(kr' \text{sen} \theta) + \left(\frac{j^{(n-1)} + (j)^{(2n)}}{j^{(n-1)}} \right) J_{(n-1)}(kr' \text{sen} \theta) \right\} = \\ & = -\frac{\pi \text{sen} n\phi}{2} \left\{ (j^{(n+1)} - j^{(n-1)}) J_{(n+1)}(kr' \text{sen} \theta) + (j^{(n-1)} - j^{(n+1)}) J_{(n-1)}(kr' \text{sen} \theta) \right\} + \\ & + \frac{\pi \cos n\phi}{2j} \left\{ (j^{(n+1)} + j^{(n-1)}) J_{(n+1)}(kr' \text{sen} \theta) + (j^{(n-1)} + j^{(n+1)}) J_{(n-1)}(kr' \text{sen} \theta) \right\} = \end{aligned}$$

OU

$$\begin{aligned} & \int_0^{2\pi} \text{sen}(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen} \theta \cos(\phi - \phi')} d\phi' = \\ & = -\frac{\pi \text{sen} n\phi}{2} (j^{(n+1)} - j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen} \theta) - J_{(n-1)}(kr' \text{sen} \theta) \} + \frac{\pi \cos n\phi}{2j} (j^{(n+1)} + j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen} \theta) + J_{(n-1)}(kr' \text{sen} \theta) \} = \\ & = -\frac{\pi \text{sen} n\phi}{2} (j^{(n+1)} - j^{(n-1)}) \{ J_{(n+1)}(kr' \text{sen} \theta) - J_{(n-1)}(kr' \text{sen} \theta) \} = -\frac{\pi \text{sen} n\phi}{2} j^{(n)} (j - j^{-1}) \{ J_{(n+1)}(kr' \text{sen} \theta) - J_{(n-1)}(kr' \text{sen} \theta) \} = \\ & = -\frac{\pi \text{sen} n\phi}{2} j^{(n)} (-2j^{-1}) \{ J_{(n+1)}(kr' \text{sen} \theta) - J_{(n-1)}(kr' \text{sen} \theta) \} = \pi \text{sen} n\phi j^{(n-1)} \{ J_{(n+1)}(kr' \text{sen} \theta) - J_{(n-1)}(kr' \text{sen} \theta) \} = \\ & = -2\pi \text{sen} n\phi j^{(n-1)} J'_{(n)}(kr' \text{sen} \theta) \end{aligned}$$

logo:

$$\int_0^{2\pi} \text{sen}(n\phi') \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = -2\pi \text{sen } n\phi j^{(n-1)} J_n'(kr' \text{sen } \theta) \quad (\text{E.4})$$

E.5

Obtenção das Integrais para o Caso Especial em que $n = 1$

Com $n=1$, a Equação (E.2) pode ser reescrita como:

$$\int_0^{2\pi} \cos(\phi') \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = 2\pi \cos \phi j^{(-1)} J_1'(kr' \text{sen } \theta)$$

mas,

$$\begin{aligned} \int_0^{2\pi} \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' &= -\int_{\phi}^{\phi-2\pi} \cos(\alpha) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = -\frac{1}{2} \int_{\phi}^{\phi-2\pi} (e^{j\alpha} + e^{-j\alpha}) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = \\ \phi - \phi' = \alpha &\Rightarrow \phi' = \phi - \alpha \\ d\phi' &= -d\alpha \\ &= -\frac{1}{2} \left(\int_{\phi}^{\phi-2\pi} e^{j\alpha} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha + \int_{\phi}^{\phi-2\pi} e^{-j\alpha} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha \right) = \\ &= -2\pi \cos \phi J_2(kr' \text{sen } \theta) \end{aligned}$$

então:

$$\begin{aligned} \int_0^{2\pi} \cos 2\phi' e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' &= \frac{1}{2} \int_0^{2\pi} (e^{j2\phi'} + e^{-j2\phi'}) e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \\ &= \frac{1}{2} \int_0^{2\pi} e^{j2\phi'} e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' + \frac{1}{2} \int_0^{2\pi} e^{-j2\phi'} e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \pi J_2(kr' \text{sen } \theta) (e^{j2\phi} + e^{-j2\phi}) = \\ &= 2\pi \cos 2\phi J_2(kr' \text{sen } \theta) \\ \frac{1}{2} \int_0^{2\pi} e^{j2\phi'} e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' &= \frac{1}{2} \int_{-\phi}^{2\pi-\phi} e^{j2\phi} e^{j(kr' \text{sen } \theta \cos \beta + 2\beta)} d\beta = -\frac{e^{j2\phi}}{2} \int_0^{2\pi} e^{j(kr' \text{sen } \theta \cos \beta + 2\beta)} d\beta = \\ &= -\frac{e^{j2\phi}}{2} \int_0^{2\pi} e^{j2\beta} e^{j(kr' \text{sen } \theta \cos \beta)} d\beta = -\frac{e^{j2\phi}}{2} \frac{2\pi}{j^{-2}} J_2(kr' \text{sen } \theta) e^{-j2\phi} = \pi e^{j2\phi} J_2(kr' \text{sen } \theta) \\ \beta = \phi' - \phi &\Rightarrow \phi' = \beta + \phi \\ d\beta &= d\phi' \\ \phi' = 0 &\Rightarrow \beta = -\phi \\ \phi' = 2\pi &\Rightarrow \beta = 2\pi - \phi \\ \frac{1}{2} \int_0^{2\pi} e^{-j2\phi'} e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' &= \frac{1}{2} \int_{-\phi}^{2\pi-\phi} e^{-j2\phi} e^{j(kr' \text{sen } \theta \cos \beta - 2\beta)} d\beta = -\frac{e^{-j2\phi}}{2} \int_0^{2\pi} e^{j(kr' \text{sen } \theta \cos \beta - 2\beta)} d\beta = \\ &= -\frac{e^{-j2\phi}}{2} \int_0^{2\pi} e^{-j2\beta} e^{j(kr' \text{sen } \theta \cos \beta)} d\beta = -\frac{e^{-j2\phi}}{2} \frac{2\pi}{j^2} J_{-2}(kr' \text{sen } \theta) e^{-j2\phi} = \pi e^{-j2\phi} J_2(kr' \text{sen } \theta) \\ \frac{j^{-n}}{2\pi} \int_0^{2\pi} e^{jn\beta} e^{jx \cos \beta} d\beta &= J_n(x) \end{aligned}$$

logo:

$$\int_0^{2\pi} \cos \phi' \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \pi \cos \phi [J_0(kr' \text{sen } \theta) - J_2(kr' \text{sen } \theta)]$$

Do mesmo modo, obtém-se:

$$\int_0^{2\pi} \cos \phi' \text{sen}(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \pi \text{sen } \phi (J_0(kr' \text{sen } \theta) + J_2(kr' \text{sen } \theta))$$

$$\int_0^{2\pi} \text{sen } \phi' \cos(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = \pi \text{sen } \phi [J_0(kr' \text{sen } \theta) - J_2(kr' \text{sen } \theta)]$$

$$\int_0^{2\pi} \text{sen } \phi' \text{sen}(\phi - \phi') e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = -\pi \cos \phi [J_0(kr' \text{sen } \theta) + J_2(kr' \text{sen } \theta)]$$

$$\int_0^{2\pi} \cos 2\alpha e^{ju \cos \alpha} d\alpha = \int_0^{2\pi} \left(\frac{e^{i2\alpha} + e^{-i2\alpha}}{2} \right) e^{ju \cos \alpha} d\alpha = -2\pi J_2(u)$$

$$\int_0^{2\pi} \text{sen } 2\alpha e^{ju \cos \alpha} d\alpha = \int_0^{2\pi} \left(\frac{e^{i2\alpha} - e^{-i2\alpha}}{2j} \right) e^{ju \cos \alpha} d\alpha = 0$$

$$\int_0^{2\pi} e^{i2\alpha} e^{ju \cos \alpha} d\alpha = -2\pi J_2(u)$$

$$\int_0^{2\pi} e^{ju \cos \alpha} d\alpha = -2\pi J_0(u)$$

De [14, (IV-35)]:

$$\int_0^{2\pi} e^{jn\alpha} e^{ju \cos \alpha} d\alpha = \frac{2\pi}{j^{-n}} J_n(u)$$

então:

$$\int_0^{2\pi} e^{jn\alpha} e^{ju \cos \alpha} d\alpha = \frac{2\pi}{j^{-n}} J_n(u)$$

$$\int_0^{2\pi} \cos n\alpha e^{ju \cos \alpha} d\alpha = \int_0^{2\pi} \left(\frac{e^{in\alpha} + e^{-in\alpha}}{2} \right) e^{ju \cos \alpha} d\alpha =$$

$$= \frac{1}{2} \left(\frac{2\pi}{j^{-n}} J_n(u) + \frac{2\pi}{j^n} J_{-n}(u) \right) = \frac{1}{2} \left(\frac{2\pi}{j^{-n}} J_n(u) + \frac{2\pi}{j^n} (-1)^n J_n(u) \right)$$

$$J_{-n}(X) = (-1)^n J_n(X)$$

$$n = 0 \Rightarrow 2\pi J_0(u)$$

$$n = 1 \Rightarrow \pi J_1(u) \left(\frac{1}{j^{-1}} - \frac{1}{j^1} \right) =$$

$$N = 2 \Rightarrow \frac{1}{2} \left(\frac{2\pi}{j^{-2}} J_2(u) + \frac{2\pi}{j^2} (-1)^2 J_2(u) \right) = -2\pi J_2(u)$$

ou

$$\int_0^{2\pi} \cos 2\phi' e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = 2\pi \cos 2\phi J_2(kr' \text{sen } \theta)$$

E.6 Desenvolvimento para os Casos Apresentados

a)

$$\int_0^{2\pi} \text{sen}^2 \phi' e^{jkr' \text{sen } \theta \cos(\phi - \phi')} d\phi' = - \int_{\phi}^{\phi - 2\pi} \text{sen}^2(\phi - \alpha) e^{jkr' \text{sen } \theta \cos \alpha} d\alpha = - \frac{1}{2} \int_{\phi}^{\phi - 2\pi} [1 - \cos(2\phi - 2\alpha)] e^{jkr' \text{sen } \theta \cos \alpha} d\alpha =$$

$$\alpha = \phi - \phi'$$

$$d\alpha = -d\phi'$$

$$= \frac{1}{2} \int_0^{2\pi} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha + \frac{1}{2} \int_{\phi}^{\phi - 2\pi} \cos 2\phi \cos 2\alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha + \frac{1}{2} \int_{\phi}^{\phi - 2\pi} \text{sen } 2\phi \text{sen } 2\alpha e^{jkr' \text{sen } \theta \cos \alpha} d\alpha =$$

$$= -\pi J_0(kr' \text{sen } \theta) + \frac{1}{2} \cos 2\phi \int_{\phi}^{\phi - 2\pi} \frac{(e^{j2\phi'} + e^{-j2\phi'})}{2} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha + \frac{1}{2} \text{sen } 2\phi \int_{\phi}^{\phi - 2\pi} \frac{(e^{j2\phi'} - e^{-j2\phi'})}{2j} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha =$$

$$= -\pi J_0(kr' \text{sen } \theta) - \frac{1}{4} \cos 2\phi \left\{ \int_0^{2\pi} e^{j2\phi'} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha + \int_0^{2\pi} e^{-j2\phi'} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha \right\} +$$

$$\frac{1}{4j} \text{sen } 2\phi \left\{ \int_0^{2\pi} e^{j2\phi'} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha - \int_0^{2\pi} e^{-j2\phi'} e^{jkr' \text{sen } \theta \cos \alpha} d\alpha \right\} =$$

$$= -\pi J_0(kr' \text{sen } \theta) - \frac{1}{4} \cos 2\phi \{ -2\pi J_2(kr' \text{sen } \theta) - 2\pi J_2(kr' \text{sen } \theta) \} =$$

$$= -\pi J_0(kr' \text{sen } \theta) + \pi \cos 2\phi J_2(kr' \text{sen } \theta)$$

b)

$$\begin{aligned}
& \int_0^{2\pi} \cos \phi' \operatorname{sen}(\phi - \phi') e^{jkr' \operatorname{sen} \theta \cos(\phi - \phi')} d\phi' = \\
& \alpha = \phi - \phi' \Rightarrow \phi' = -(\alpha - \phi) \\
& d\alpha = -d\phi' \\
& = -\int_{\phi}^{\phi-2\pi} \cos(\alpha - \phi) \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = -\int_{\phi}^{\phi-2\pi} \cos \phi \cos \alpha \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \int_{\phi}^{\phi-2\pi} \operatorname{sen} \phi \operatorname{sen}^2 \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
& = -\frac{\cos \phi}{2} \int_{\phi}^{\phi-2\pi} \operatorname{sen} 2\alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \frac{\operatorname{sen} \phi}{2} \int_{\phi}^{\phi-2\pi} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha + \frac{\operatorname{sen} \phi}{2} \int_{\phi}^{\phi-2\pi} \cos 2\alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
& = -\frac{\cos \phi}{2} \int_0^{2\pi} \frac{(e^{j2\phi'} - e^{-j2\phi'})}{2j} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha - \frac{\operatorname{sen} \phi}{2} (-2\pi J_0(kr' \operatorname{sen} \theta)) - \frac{\operatorname{sen} \phi}{2} \int_0^{2\pi} \frac{(e^{j2\phi'} + e^{-j2\phi'})}{2} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha \\
& = -\frac{\cos \phi}{4j} \{-2\pi J_2(kr' \operatorname{sen} \theta) + 2\pi J_2(kr' \operatorname{sen} \theta)\} + \operatorname{sen} \phi \pi J_0(kr' \operatorname{sen} \theta) - \frac{\operatorname{sen} \phi}{4} \{-2\pi J_2(kr' \operatorname{sen} \theta) - 2\pi J_2(kr' \operatorname{sen} \theta)\} = \\
& = \pi \operatorname{sen} \phi [J_0(kr' \operatorname{sen} \theta) + J_2(kr' \operatorname{sen} \theta)]
\end{aligned}$$

c)

$$\begin{aligned}
& \int_0^{2\pi} \operatorname{sen} \phi' \operatorname{sen}(\phi - \phi') e^{jkr' \operatorname{sen} \theta \cos(\phi - \phi')} d\phi' = ? \\
& \alpha = \phi - \phi' \Rightarrow \phi' = -(\alpha - \phi) \\
& d\alpha = -d\phi' \\
& = -\int_{\phi}^{\phi-2\pi} \operatorname{sen}(\alpha - \phi) \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = -\int_{\phi}^{\phi-2\pi} \cos \phi \operatorname{sen}^2 \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha + \int_{\phi}^{\phi-2\pi} \operatorname{sen} \phi \cos \alpha \operatorname{sen} \alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
& = -\frac{\cos \phi}{2} \int_{\phi}^{\phi-2\pi} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha + \frac{\cos \phi}{2} \int_{\phi}^{\phi-2\pi} \cos 2\alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha + \frac{\operatorname{sen} \phi}{2} \int_{\phi}^{\phi-2\pi} \operatorname{sen} 2\alpha e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
& = -\frac{\cos \phi}{2} (2\pi J_0(kr' \operatorname{sen} \theta)) + \frac{\cos \phi}{2} \int_0^{2\pi} \frac{(e^{j2\phi'} + e^{-j2\phi'})}{2} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha + \frac{\operatorname{sen} \phi}{2} \int_0^{2\pi} \frac{(e^{j2\phi'} - e^{-j2\phi'})}{2j} e^{jkr' \operatorname{sen} \theta \cos \alpha} d\alpha = \\
& = -\cos \phi \pi J_0(kr' \operatorname{sen} \theta) + \frac{\cos \phi}{4} \{-2\pi J_2(kr' \operatorname{sen} \theta) - 2\pi J_2(kr' \operatorname{sen} \theta)\} + \frac{\operatorname{sen} \phi}{4} \{-2\pi J_2(kr' \operatorname{sen} \theta) + 2\pi J_2(kr' \operatorname{sen} \theta)\} = \\
& = -\cos \phi \pi J_0(kr' \operatorname{sen} \theta) - \frac{4\pi \cos \phi}{4} \{J_2(kr' \operatorname{sen} \theta)\} = -\pi \cos \phi J_0(kr' \operatorname{sen} \theta) - \pi \cos \phi J_2(kr' \operatorname{sen} \theta) = \\
& = -\pi \cos \phi [J_0(kr' \operatorname{sen} \theta) + J_2(kr' \operatorname{sen} \theta)]
\end{aligned}$$