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A

Integrals of Lipschitz-Hankel type involving products of Bessel functions

The integrals of Lipschitz-Hankel type involving products of Bessel functions can be represented by

$$I_{pq\lambda}(\xi, r; c) = \int_0^\infty J_p(\xi t) J_q(rt) e^{-ct} t^\lambda dt \quad (\text{A-1})$$

where p , q and λ are integers; and $J_p(\xi t)$ and $J_q(rt)$ are Bessel functions of the first kind of order p and q , respectively. The convergent integrals of this type were tabulated by Eason et al. [67]. The expressions used in the previous chapters are

$$I_{000} = \frac{2k}{\pi A_1} K(m) \quad (\text{A-2})$$

$$I_{110} = -\frac{2(k^2 - 2)}{\pi k A_1} K(m) - \frac{4}{\pi k A_1} E(m) \quad (\text{A-3})$$

$$I_{100} = \begin{cases} -\frac{kc}{2\pi\xi\sqrt{\xi}r} K(m) - \frac{\Lambda_0(n,m)}{2\xi} + \frac{1}{\xi} & \text{if } \xi > r \\ -\frac{kc}{2\pi\xi^2} K(m) + \frac{1}{2\xi} & \text{if } \xi = r \\ -\frac{kc}{2\pi\xi\sqrt{\xi}r} K(m) + \frac{\Lambda_0(n,m)}{2\xi} & \text{if } \xi < r \end{cases} \quad (\text{A-4})$$

$$I_{001} = \frac{2ck^3}{\pi\bar{k}^2 A_1^3} E(m) \quad (\text{A-5})$$

$$I_{111} = -\frac{4ck}{\pi A_1^3} K(m) - \frac{2ck(k^2 - 2)}{\pi\bar{k}^2 A_1^3} E(m) \quad (\text{A-6})$$

$$I_{101} = \frac{k}{\pi\xi A_1} K(m) + \frac{k^3 A_2}{\pi\xi\bar{k}^2 A_1^3} E(m) \quad (\text{A-7})$$

$$I_{002} = -\frac{2c^2 k^5}{\pi\bar{k}^2 A_1^5} K(m) - \frac{2k^3}{\pi\bar{k}^2 A_1^3} \left[1 + \frac{2c^2 k^2 (k^2 - 2)}{\bar{k}^2 A_1^2} \right] E(m) \quad (\text{A-8})$$

$$I_{112} = \frac{2k}{\pi A_1^3} \left[2 + \frac{c^2 k^2 (k^2 - 2)}{\bar{k}^2 A_1^2} \right] K(m) + \frac{2k}{\pi\bar{k}^2 A_1^3} \left[k^2 - 2 + \frac{2c^2 k^2 (k^4 + \bar{k}^2)}{\bar{k}^2 A_1^2} \right] E(m) \quad (\text{A-9})$$

$$I_{102} = -\frac{ck^5 A_2}{\pi\xi\bar{k}^2 A_1^5} K(m) + \frac{ck^3}{\pi\xi\bar{k}^2 A_1^3} \left[3 - \frac{2k^2 A_2 (k^2 - 2)}{\bar{k}^2 A_1^2} \right] E(m) \quad (\text{A-10})$$

$$I_{003} = \frac{2 c k^5}{\pi \bar{k}^2 A_1^5} \left[3 + \frac{4 c^2 k^2 (k^2 - 2)}{\bar{k}^2 A_1^2} \right] K(m) - \frac{2 c k^5}{\pi \bar{k}^4 A_1^5} \left[-6(k^2 - 2) - \frac{c^2 k^2}{A_1^2} \left(\frac{8 k^4}{\bar{k}^2} + 23 \right) \right] E(m) \quad (\text{A-11})$$

$$I_{113} = - \frac{2 c k^3}{\pi \bar{k}^2 A_1^5} \left[3(k^2 - 2) + \frac{2 c^2 k^2}{A_1^2} \left(\frac{2 k^4}{\bar{k}^2} + 3 \right) \right] K(m) - \frac{2 c k^3}{\pi \bar{k}^4 A_1^5} \left[6(k^4 + \bar{k}^2) + \frac{c^2 k^2 (k^2 - 2)(3\bar{k}^2 + 8k^4)}{\bar{k}^2 A_1^2} \right] E(m) \quad (\text{A-12})$$

$$I_{103} = \frac{k^5}{\pi \xi \bar{k}^2 A_1^5} \left[A_2 - 5c^2 + \frac{4 k^2 (k^2 - 2) c^2 A_2}{\bar{k}^2 A_1^2} \right] K(m) + \frac{k^3}{\pi \xi \bar{k}^2 A_1^3} \left[-3 + \frac{2 k^2 (k^2 - 2) (A_2 - 5c^2)}{\bar{k}^2 A_1^2} + \frac{c^2 k^4 A_2}{\bar{k}^2 A_1^4} \left(\frac{8 k^4}{\bar{k}^2} + 23 \right) \right] E(m) \quad (\text{A-13})$$

$$I_{004} = - \frac{2 k^5}{\pi \bar{k}^2 A_1^5} \left[3 + \frac{24 c^2 k^2 (k^2 - 2)}{\bar{k}^2 A_1^2} + \frac{c^4 k^4 (24k^4 + 41\bar{k}^2)}{\bar{k}^4 A_1^4} \right] K(m) - \frac{4 k^5}{\pi \bar{k}^4 A_1^5} \left[3(k^2 - 2) + \frac{3 c^2 k^2 (8k^4 + 23\bar{k}^2)}{\bar{k}^2 A_1^2} + \frac{4 c^4 k^4 (k^2 - 2)(6k^4 + 11\bar{k}^2)}{\bar{k}^4 A_1^4} \right] E(m) \quad (\text{A-14})$$

$$I_{114} = \frac{6 k^3}{\pi \bar{k}^2 A_1^5} \left[k^2 - 2 + \frac{4 c^2 k^2 (2k^4 + 3\bar{k}^2)}{\bar{k}^2 A_1^2} + \frac{c^4 k^4 (k^2 - 2)(8k^4 + 5\bar{k}^2)}{\bar{k}^4 A_1^4} \right] K(m) + \frac{12 k^3}{\pi \bar{k}^4 A_1^5} \left[k^2 + \bar{k}^4 + \frac{c^2 k^2 (k^2 - 2)(3\bar{k}^2 + 8k^4)}{\bar{k}^2 A_1^2} + \frac{c^4 k^4 (8 - 4\bar{k}^2 - 3\bar{k}^4 - 4\bar{k}^6 + 8\bar{k}^8)}{\bar{k}^4 A_1^4} \right] E(m) \quad (\text{A-15})$$

$$I_{104} = - \frac{c k^5}{\pi \xi \bar{k}^2 A_1^5} \left[-15 + \frac{4 k^2 (k^2 - 2)(3A_2 - 7c^2)}{\bar{k}^2 A_1^2} + \frac{c^2 k^4 A_2 (24k^4 + 71\bar{k}^2)}{\bar{k}^4 A_1^4} \right] K(m) - \frac{c k^5}{\pi \xi \bar{k}^4 A_1^5} \left[-30(k^2 - 2) + \frac{k^2 (8k^4 + 23\bar{k}^2)(3A_2 - 7c^2)}{\bar{k}^2 A_1^2} + \frac{8c^3 k^4 A_2 (k^2 - 2)(6k^4 + 11\bar{k}^2)}{\bar{k}^4 A_1^4} \right] E(m) \quad (\text{A-16})$$

in which $I_{pq\lambda}(\xi, r; c) = I_{qp\lambda}(r, \xi; c)$ and

$$A_1 = 2 \sqrt{\xi r}, \quad A_2 = \xi^2 - r^2 - c^2, \quad A_3 = -\xi^2 + r^2 - c^2 \quad (\text{A-17})$$

In the above expressions, $K(m)$ and $E(m)$ are the complete elliptic integrals of the first and second kinds, respectively,

$$K(m) = \int_0^{\pi/2} (1 - m \sin \theta^2)^{-1/2} d\theta \quad (\text{A-18})$$

$$E(m) = \int_0^{\pi/2} (1 - m \sin \theta^2)^{1/2} d\theta \quad (\text{A-19})$$

the modulus k , the complementary modulus \bar{k} and the parameter m are given by

$$k = \frac{2 \sqrt{\xi r}}{\sqrt{(\xi + r)^2 + c^2}}, \quad \bar{k} = \sqrt{1 - k^2} \quad \text{and} \quad m = k^2 \quad (\text{A-20})$$

In Eq. (A-4), $\Lambda_0(n, m)$ is the Heuman complete elliptic integral expressed as

$$\Lambda_0(n, m) = \frac{2}{\pi} \left[\sqrt{1 - n} \sqrt{1 - \frac{m}{n}} \Pi(n, m) \right] \quad (\text{A-21})$$

where $\Pi(n, m)$ is the complete elliptic integral of the third kind defined as

$$\Pi(n, m) = \int_0^{\pi/2} (1 - n \sin \theta^2)^{-1} (1 - m \sin \theta^2)^{-1/2} d\theta \quad (\text{A-22})$$

and n is the characteristic number

$$n = \frac{A_1^2}{(\xi + r)^2} \quad (\text{A-23})$$

Notice that all Lipschitz-Hankel integrals $I_{pq\ell}(\xi, r; c)$ listed above are written in terms of $K(m)$, $E(m)$ and $\Pi(m)$, which can be numerically evaluated by duplication as proposed by Carlson [93, 94].

In the following, some useful limits are given, for $\rho_0 = \sqrt{r^2 + c^2}$:

$$\lim_{\xi \rightarrow 0} I_{000} = \frac{1}{\rho_0} \quad (\text{A-24})$$

$$\lim_{\xi \rightarrow 0} I_{110}/\xi = \frac{r}{2\rho_0^3} \quad (\text{A-25})$$

$$\lim_{\xi \rightarrow 0} I_{100}/\xi = \frac{c}{2\rho_0^3} \quad (\text{A-26})$$

$$\lim_{\xi \rightarrow 0} I_{010} = \frac{c - \rho_0}{r\rho_0} \quad (\text{A-27})$$

$$\lim_{\xi \rightarrow 0} I_{001} = \frac{c}{\rho_0^3} \quad (\text{A-28})$$

$$\lim_{\xi \rightarrow 0} I_{111}/\xi = \frac{3rc}{2\rho_0^5} \quad (\text{A-29})$$

$$\lim_{\xi \rightarrow 0} I_{101}/\xi = -\frac{r^2 - 2c^2}{2\rho_0^5} \quad (\text{A-30})$$

$$\lim_{\xi \rightarrow 0} I_{011} = \frac{r}{\rho_0^3} \quad (\text{A-31})$$

$$\lim_{\xi \rightarrow 0} I_{002} = -\frac{r^2 - 2c^2}{\rho_0^5} \quad (\text{A-32})$$

$$\lim_{\xi \rightarrow 0} I_{112}/\xi = -\frac{3r(r^2 - 4c^2)}{2\rho_0^7} \quad (\text{A-33})$$

$$\lim_{\xi \rightarrow 0} I_{102}/\xi = -\frac{3c(3r^2 - 2c^2)}{2\rho_0^7} \quad (\text{A-34})$$

$$\lim_{\xi \rightarrow 0} I_{012} = \frac{3rc}{\rho_0^5} \quad (\text{A-35})$$

$$\lim_{\xi \rightarrow 0} I_{003} = -\frac{3c(3r^2 - 2c^2)}{\rho_0^7} \quad (\text{A-36})$$

$$\lim_{\xi \rightarrow 0} I_{113}/\xi = -\frac{15rc(3r^2 - 4c^2)}{2\rho_0^9} \quad (\text{A-37})$$

$$\lim_{\xi \rightarrow 0} I_{103}/\xi = \frac{3(3r^4 - 24r^2 c^2 + 8c^4)}{2\rho_0^9} \quad (\text{A-38})$$

$$\lim_{\xi \rightarrow 0} I_{114}/\xi = \frac{45r(r^4 - 12r^2 c^2 + 8c^4)}{2\rho_0^{11}} \quad (\text{A-41})$$

$$\lim_{\xi \rightarrow 0} I_{013} = -\frac{3r(r^2 - 4c^2)}{\rho_0^7} \quad (\text{A-39})$$

$$\lim_{\xi \rightarrow 0} I_{104}/\xi = \frac{15c(15r^4 - 40r^2 c^2 + 8c^4)}{2\rho_0^{11}} \quad (\text{A-42})$$

$$\lim_{\xi \rightarrow 0} I_{004} = \frac{3(3r^4 - 24r^2 c^2 + 8c^4)}{\rho_0^9} \quad (\text{A-40})$$

$$\lim_{\xi \rightarrow 0} I_{014} = -\frac{15cr(3r^2 - 4c^2)}{\rho_0^9} \quad (\text{A-43})$$

The expressions listed in Eqs. (A-2) to (A-16) and Eqs. (A-24) to (A-43) are part of the integrand of the boundary integrals arising in the formulations presented in this work. Depending on the value of m and ρ_0 , some of these integrals become weakly and others strongly singular, as dealt with in Appendix B.

B

Numerical integration

This appendix presents the numerical schemes used to evaluate the integrals arising in the boundary element formulations for axisymmetric problems. As only the meridian of the axisymmetric boundary needs to be discretized, these integrals are calculated along the boundary $\Gamma(r, z)$, for each portion between consecutive nodes of an element.

B.1 Regular integral

Let $f(r, z)$ be a regular function on Γ , in the sense that it can be approximated by a polynomial of a not too high degree in the domain of interest. Then, its integral can be expressed in a natural coordinate system η in the interval $[-1, 1]$ and approximated by the Gauss-Legendre quadrature rule [86], arriving at

$$\int_{\Gamma} f(r, z) d\Gamma = \int_{-1}^1 f(\eta) J(\eta) d\eta \cong \sum_{m=1}^{n_g} [f(\eta) J(\eta)]_{\eta=\eta_m^g} w_m^g \quad (\text{B-1})$$

where

$$J(\eta) = \sqrt{\left(\frac{dr}{d\eta}\right)^2 + \left(\frac{dz}{d\eta}\right)^2} \quad (\text{B-2})$$

is the Jacobian transformation between the global and natural coordinate systems. The coefficients η_m^g and w_m^g are the abscissas and weights of the Gauss-Legendre quadrature rule for n_g points within the interval $(-1, 1)$, which suffice to exactly evaluate the integral of a polynomial of order $2n_g - 1$.

B.2 Weakly singular integral of logarithmic terms

Let $f(r, z)$ be a regular function and $\rho(r, z)$ the distance between the points $P(\xi, z')$ and $Q(r, z)$ on the boundary $\Gamma(r, z)$. One is concerned with the evaluation of the following weakly singular integral

$$\int_{\Gamma} f(r, z) \ln \rho(r, z) d\Gamma = \int_{-1}^1 f(\eta) \ln \rho(\eta) J(\eta) d\eta \quad (\text{B-3})$$

for the case $\rho(-1) = 0$ or $\rho(1) = 0$. A unified treatment of both cases may be obtained by expressing

$$\rho(\eta) = \bar{\rho}(\eta, \eta')(1 - \eta' \eta) \quad (\text{B-4})$$

where η' is equal to either -1 or 1 and $\bar{\rho}(\eta, \eta')$ is the non-vanishing part of $\rho(\eta)$ for $\eta \in (-1, 1)$. Then, the integral of Eq. (B-3) may be decomposed as [74]

$$\int_{\Gamma} f(r, z) \ln \rho(r, z) d\Gamma = \int_{-1}^1 f(\eta) \ln[2\bar{\rho}(\eta)] J(\eta) d\eta + 2 \int_0^1 f(\tilde{\eta}) \ln \tilde{\eta} J(\tilde{\eta}) d\tilde{\eta} \quad (\text{B-5})$$

in which the transformation to the natural coordinate system $\tilde{\eta} \in [0, 1]$ is given by

$$\tilde{\eta} = \frac{1}{2}(1 - \eta' \eta) \quad (\text{B-6})$$

The resulting integrals can be approximated by the Gauss-Legendre and logarithmic weighted Gauss quadratures rules [86], leading to

$$\begin{aligned} \int_{\Gamma} f(r, z) \ln \rho(r, z) d\Gamma &\cong \sum_{m=1}^{n_g} [f(\eta) \ln[2\bar{\rho}(\eta)] J(\eta)]_{\eta=\eta_m^g} w_m^g + \\ &\quad \sum_{m=1}^{n_l} [f(\tilde{\eta}) \ln \tilde{\eta} J(\tilde{\eta})]_{\tilde{\eta}=\eta_m^l} w_m^l \end{aligned} \quad (\text{B-7})$$

The coefficients η_m^l and w_m^l are the abscissas and weights of the logarithmic weighted Gauss quadrature rule for n_l points within the interval (0, 1), which suffice to exactly evaluate the integral of a polynomial of order $n_l - 1$.

The above integration scheme is obtained from a transformation of variables and the use of Gauss-Legendre and logarithmic weighted Gauss quadratures rules. However, other approaches can also be employed, such as performing a transformation of variables so that the singular terms vanish, or regularizing the kernels to be evaluated in the numerical scheme with analytical integration of the regular part [74].

B.2.1

Weakly singular integral of terms with the complete elliptic integral of the first order

Let $f(r, z)$ be a regular function and $K(m)$ the complete elliptic integral of the first order given by Eq. (A-18), of modulus

$$m = \frac{4\xi r}{(\xi + r)^2 + (z' - z)^2} \quad (\text{B-8})$$

given in terms of the coordinates of points P(ξ, z') and Q(r, z) on the boundary $\Gamma(r, z)$. One is concerned with the evaluation of the following weakly singular integral

$$\int_{\Gamma} f(r, z) K(m) d\Gamma = \int_{-1}^1 f(\eta) K(m) J(\eta) d\eta \quad (\text{B-9})$$

which actually encompasses to singularity in the case of $K(m) = \infty$ since $m = 1$ for either $\eta = -1$ or $\eta = 1$. The integration scheme to be presented was proposed by Bialecki et al. [87], by approximating the complete elliptic integral $K(m)$ and isolating its singular term.

The complete elliptic integral $K(m)$ can be approximated, for $0 \leq m < 1$ and within an error $\epsilon < 2 \cdot 10^{-8}$, by the expression [68]

$$K(m) = K_1(\bar{m}) - K_2(\bar{m}) \ln \bar{m} \quad (\text{B-10})$$

where

$$\bar{m} = \frac{\rho^2}{(\xi + r)^2 + (z' - z)^2} \quad (\text{B-11})$$

is the complementary modulus of the complete elliptic integral and

$$\begin{aligned} K_1(\bar{m}) &= a_0 + a_1 \bar{m} + \dots + a_4 \bar{m}^4 \\ K_2(\bar{m}) &= b_0 + b_1 \bar{m} + \dots + b_4 \bar{m}^4 \end{aligned} \quad (\text{B-12})$$

are polynomials whose coefficients are given by

$$\begin{aligned} a_0 &= 1.38629436112 & b_0 &= 0.5 \\ a_1 &= 0.09666344259 & b_1 &= 0.12498593597 \\ a_2 &= 0.03590092383 & b_2 &= 0.06880248576 \\ a_3 &= 0.03742563713 & b_3 &= 0.03328355346 \\ a_4 &= 0.01451196212 & b_4 &= 0.00441787012 \end{aligned} \quad (\text{B-13})$$

Substituting the approximation given by Eq. (B-10) into the weakly singular integral in Eq. (B-9), one may isolate the singular term to obtain

$$\begin{aligned} \int_{\Gamma} f(r, z) K(m) d\Gamma &= \int_{\Gamma} f(r, z) \left[K_1(\bar{m}) + K_2(\bar{m}) \ln \frac{\rho(r, z)^2}{\bar{m}} \right] d\Gamma - \\ &\quad 2 \int_{\Gamma} f(r, z) K_2(\bar{m}) \ln \rho(r, z) d\Gamma \end{aligned} \quad (\text{B-14})$$

Applying the scheme for the regular integral presented in Section B.1 to the first integral and the scheme for the weakly singular integral in Section B.2 to the second integral, one obtains

$$\begin{aligned} \int_{\Gamma} f(r, z) K(m) d\Gamma &\cong \sum_{m=1}^{n_g} \left\{ f(\eta) \left[K_1(\bar{m}) + 2K_2(\bar{m}) \ln \frac{1 - \eta' \eta}{2 \sqrt{\bar{m}}} \right] J(\eta) \right\}_{\eta=\eta_m^s} w_m^s - \\ &\quad 4 \sum_{m=1}^{n_l} [f(\tilde{\eta}) K_2(\bar{m}) \ln \tilde{\eta} J(\tilde{\eta})]_{\tilde{\eta}=\eta_m^l} w_m^l \end{aligned} \quad (\text{B-15})$$

for $\tilde{\eta}$ according to the development that has led to Eq. (B-6).

B.2.2

Weakly singular integral of terms with the complete elliptic integral of the second order

Let $f(r, z)$ be a regular function and $E(m)$ the complete elliptic integral of the second order, of modulus m , given by Eq. (A-19) in terms of the coordinates of points P(ξ, z') and Q(r, z) on the boundary $\Gamma(r, z)$. One needs to evaluate the following weakly singular integral

$$\int_{\Gamma} f(r, z) E(m) d\Gamma = \int_{-1}^1 f(\eta) E(m) J(\eta) d\eta \quad (\text{B-16})$$

for the case of $m = 1$ for either $\eta = -1$ or $\eta = 1$. Although $E(m) \neq \infty$ for this case, one may isolate the quasi-singular terms to enhance the convergence of the numerical integration.

The complete elliptic integral $E(m)$ can be approximated, for $0 \leq m < 1$ and within an error $\epsilon < 2 \cdot 10^{-8}$, by the expression [68]

$$E(m) = E_1(\bar{m}) - E_2(\bar{m}) \ln \bar{m} \quad (\text{B-17})$$

where \bar{m} is given by Eq. (B-8) and

$$\begin{aligned} E_1(\bar{m}) &= 1 + a_1 \bar{m} + \dots + a_4 \bar{m}^4 \\ E_2(\bar{m}) &= b_1 \bar{m} + \dots + b_4 \bar{m}^4 \end{aligned} \quad (\text{B-18})$$

are the polynomials whose coefficients are given by

$$\begin{aligned} a_1 &= 0.44325141463 & b_1 &= 0.24998368310 \\ a_2 &= 0.06260601220 & b_2 &= 0.09200180037 \\ a_3 &= 0.04757383546 & b_3 &= 0.04069697526 \\ a_4 &= 0.01736506451 & b_4 &= 0.00526449639 \end{aligned} \quad (\text{B-19})$$

The polynomial approximation of $E(m)$ presents no singularity, since $E_2(\bar{m})$ has no free coefficients, according to Eq. (B-17). However, the presence of $\ln \bar{m}$ causes the integrand of Eq. (B-16) to be non-analytical, which requires a special numerical treatment.

In a manner similar to that used in the previous section, one arrives at the following expression for the numerical evaluation of the weakly singular integral given by Eq. (B-16)

$$\begin{aligned} \int_{\Gamma} f(r, z) E(m) d\Gamma &\cong \sum_{m=1}^{n_g} \left\{ f(\eta) \left[E_1(\bar{m}) + 2E_2(\bar{m}) \ln \frac{1 - \eta' \eta}{2 \sqrt{\bar{m}}} \right] J(\eta) \right\}_{\eta=\eta_m^g} w_m^g - \\ &4 \sum_{m=1}^{n_l} [f(\tilde{\eta}) E_2(\bar{m}) \ln \tilde{\eta} J(\tilde{\eta})]_{\tilde{\eta}=\eta_m^l} w_m^l \end{aligned} \quad (\text{B-20})$$

for $\tilde{\eta}$ given by Eq. (B-6).

B.3

Cauchy principal value of the singular integral of order $1/\rho$

Let $f(r, z)$ be a regular function and $\rho(r, z)$ the distance between the points $P(\xi, z')$ and $Q(r, z)$ on the boundary $\Gamma(r, z)$. One needs to evaluate the strongly singular integral

$$\int_{\Gamma} \frac{f(r, z)}{\rho(r, z)} d\Gamma \quad (\text{B-21})$$

for the case $\rho(-1) = 0$ or $\rho(1) = 0$. This integral may be obtained as a sum of a Cauchy principal value and a discontinuous term as

$$\int_{\Gamma} \frac{f(r, z)}{\rho(r, z)} d\Gamma = \text{PV} \int_{\Gamma} \frac{f(r, z)}{\rho(r, z)} d\Gamma + c \quad (\text{B-22})$$

The evaluation of the discontinuous term c of the strongly singular integrals appearing in the boundary element formulations is addressed in Section 3.1.5.

The Cauchy principal value is best evaluated in terms of two finite-part integrals, denoted by \oint , for the boundary segments adjacent to the singularity point $\rho(r, z) = 0$.

In what follows, one makes use of the integration scheme proposed by Dumont & Souza [88]. Using the same notation of Eq. (B-6), one may expand the regular function by Taylor and obtain the following normalized integral of Eq. (B-21) over the curved boundary Γ

$$\oint_{\Gamma} \frac{f(r, z)}{\rho(r, z)} d\Gamma = -\eta' \left[f(\eta) \ln |\bar{\rho}| \right]_{\eta=\eta'} + \oint_{-1}^1 \frac{f(\eta)}{\rho(\eta)} J(\eta) d\eta \quad (\text{B-23})$$

The resulting quadrature rule for evaluating Cauchy's principal value of the strongly singular integral of (B-21) is given by

$$\begin{aligned} \oint_{\Gamma} \frac{f(r, z)}{\rho(r, z)} d\Gamma &\cong \sum_{m=1}^{n_g} \left[\frac{f(\eta)}{\rho(\eta)} J(\eta) \right]_{\eta=\eta_m^g} w_m^g - \\ &\quad \eta' [f(\eta)]_{\eta=\eta'} \left\{ [\ln |2\bar{\rho}|]_{\eta=\eta'} - \sum_{m=1}^{n_g} \frac{w_m^g}{1 - \eta_m^g} \right\} \end{aligned} \quad (\text{B-24})$$

where

$$[\bar{\rho}(\eta)]_{\eta=\eta'} = [J(\eta)]_{\eta=\eta'} \quad (\text{B-25})$$

The above scheme, that employs the Gauss-Legendre quadrature rule and an additional correction term, evaluates exactly this integral for a polynomial function of order $2n_g - 2$. Other possible numerical integration scheme for the strongly singular integral of Eq. (B-21) makes use of Kutt's weighted quadrature rule [95] or the technique of subtraction of the singularity with further numerical integration [74]. The last is another way of expressing the procedure outlined above.